

1 DATA STRUCTURES FOR ROBUST MULTIFREQUENCY IMAGING

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4 **Abstract.** In this paper we consider imaging problems that can be cast in the form of an
5 underdetermined linear system of equations. When a single measurement vector is available, a sparsity
6 promoting ℓ_1 -minimization based algorithm may be used to solve the imaging problem efficiently. A
7 suitable algorithm in the case of multiple measurement vectors would be the MUltiple Signal Clas-
8 sification (MUSIC) which is a subspace projection method. We provide in this work a theoretical
9 framework in an abstract linear algebra setting that allows us to examine under what conditions the
10 ℓ_1 -minimization problem and the MUSIC method admit an exact solution. We also examine the
11 performance of these two approaches when the data are noisy. Several imaging configurations that
12 fall under the assumptions of the theory are discussed such as active imaging with single or multiple
13 frequency data. We also show that the phase retrieval problem can be re-cast under the same linear
14 system formalism using the polarization identity and relying on diversity of illuminations. The rele-
15 vance of our theoretical analysis in imaging is illustrated with numerical simulations and robustness
16 to noise is examined by allowing the background medium to be weakly inhomogeneous.

17 **Key words.** array imaging, phase retrieval, ℓ_1 -minimization, MUSIC

18 **1. Introduction.** Imaging is an inverse problem in which we seek to reconstruct
19 a medium's characteristics, such as the reflectivity, by recording its response to one
20 or more known excitations. The output is usually an image giving an estimate of
21 an unknown characteristic in a bounded domain, the imaging window of interest.
22 Although this problem is in all generality non-linear, it is often adequately formulated
23 as a linear system of the form

$$24 \quad (1.1) \quad \mathcal{A}\boldsymbol{\rho} = \mathbf{b},$$

25 where the data vector $\mathbf{b} \in \mathbb{C}^N$ is a linear transformation of the unknown vector $\boldsymbol{\rho} \in \mathbb{C}^K$
26 [13]. $\mathcal{A} \in \mathbb{C}^{N \times K}$ is the model matrix that relates \mathbf{b} to $\boldsymbol{\rho}$. Typically, the linear system
27 (1.1) is underdetermined because the number of unknowns K is much larger than the
28 number of measurements N , so $N \ll K$.

29 We are interested in this work in imaging problems where the unknown $\boldsymbol{\rho}$ is M-
30 sparse with $M \ll K$. Under this assumption (1.1) falls under the compressive sensing
31 framework [21, 16, 22]. It follows from [16] that the unique M-sparse solution of (1.1)
32 can be obtained with ℓ_1 -optimization when the mutual coherence¹ of the model matrix
33 \mathcal{A} is smaller than $1/(2M)$. The same result can be obtained assuming \mathcal{A} obeys the
34 M-restricted isometry property [7] which basically states that all sets of M-columns
35 of \mathcal{A} behave approximately as an orthonormal system.

We show that uniqueness for the minimal ℓ_1 solution of (1.1) can be obtained
under less restrictive conditions on the model matrix \mathcal{A} provided that the unknown $\boldsymbol{\rho}$
is such that the columns of \mathcal{A} that correspond to the support T of $\boldsymbol{\rho}$ are approximately

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¹The mutual coherence of \mathcal{A} is defined as $\max_{i \neq j} |\langle \mathbf{a}_i, \mathbf{a}_j \rangle|$ with $\mathbf{a}_i \in \mathbb{C}^N$ the columns of \mathcal{A}
normalized to one, so that $\|\mathbf{a}_i\|_{\ell_2} = 1 \forall i = 1, \dots, K$.

orthogonal, so there exists a small value $0 < \varepsilon < 1/2$ such that

$$|\langle \mathbf{a}_i, \mathbf{a}_j \rangle| < \frac{\varepsilon}{M}, \quad \forall i, j \in T, i \neq j.$$

Under this assumption, we associate to each column vector \mathbf{a}_j , $j \in T$, its vicinity

$$S_j = \left\{ k \neq j \text{ s.t. } |\langle \mathbf{a}_k, \mathbf{a}_j \rangle| \geq \frac{1}{2M} \right\}$$

36 that contains all columns of \mathcal{A} that are approximately parallel to \mathbf{a}_j . This result
 37 finds interesting applications in imaging since it states under what conditions the
 38 location of well separated reflectors can be determined with high precision. It can be
 39 also used to explain super-resolution, i.e., the significantly superior resolution that ℓ_1 -
 40 optimization provides compared to the conventional resolution of the imaging system,
 41 i.e., the Rayleigh resolution. Moreover, we address the robustness to noise of the
 42 minimal ℓ_1 solution and show that for noisy data the solution $\boldsymbol{\rho}$ can be decomposed
 43 in two parts: the coherent part $\boldsymbol{\rho}_c$, which is supported in T or in the vicinities S_j ,
 44 and the incoherent part $\boldsymbol{\rho}_i$, usually referred to as grass, that is small. Other stability
 45 results can be found in [7, 8, 17, 35, 18, 4].

46 The notion of vicinities and weak interaction between scatterers has been con-
 47 sidered in [18] and [4]. In [18], several algorithms for imaging well separated sources
 48 were introduced and analyzed. These algorithms address the issue of high coherence
 49 in \mathcal{A} using techniques of band exclusion and local optimization. In [4], a resolution
 50 analysis for ℓ_1 -minimization and ℓ_1 -penalty was carried out for array imaging in the
 51 paraxial regime. It was shown that for well separated sources or clusters of sources the
 52 minimal ℓ_1 solution is supported mainly in the vicinities of the true sources' locations.

53 More recently in [5], the problem of imaging sources in weakly inhomogeneous
 54 media was addressed using Coherent INTERferometry (CINT) followed by ℓ_1 convex
 55 optimization for deblurring. This is a natural idea since, as it was shown in [1] (see
 56 also [3]), the CINT image is a convolution of the reflectivity with a Gaussian kernel.
 57 Hence, the resolution in CINT images can be refined by deblurring as in [2], where a
 58 level set method was used. In [5], deblurring was performed with ℓ_1 -optimization and
 59 its performance was analyzed for well separated sources and well separated clusters
 60 of sources.

61 We also consider in this paper the more general form that system (1.1) takes when
 62 S multiple measurement vectors (MMV) are available, so

$$63 \quad (1.2) \quad \mathcal{A}_{l_q} \boldsymbol{\rho} = \mathbf{b}_{l_q}, \quad q = 1, \dots, S.$$

64 Here, $\mathbf{l}_q = [l_{1q}, l_{2q}, \dots, l_{Kq}]^T$ denotes a parameter vector such as the excitation that
 65 we control. To simplify the notation, we will denote the different excitations by the
 66 scalar q and write $\mathcal{A}_q \boldsymbol{\rho} = \mathbf{b}_q$ instead, unless it is necessary to explicitly state that the
 67 model matrix depends on a vector \mathbf{l}_q . To solve (1.2) we consider the Multiple Signal
 68 Classification algorithm [34] which has been used successfully in signal processing [23]
 69 and imaging [15, 25]. For a careful analysis of MUSIC for single snapshot spectral
 70 imaging we refer the reader to [26]. We show here that MUSIC gives the exact support
 71 of the solution of (1.2) in the noise free case when the matrices \mathcal{A}_q admit the following
 72 factorization

$$73 \quad (1.3) \quad \mathcal{A}_q = \tilde{\mathcal{A}} \Lambda_q, \text{ with } \Lambda_q \text{ diagonal.}$$

In this case, (1.2) admits the following MMV formulation

$$\tilde{\mathcal{A}}\boldsymbol{\rho}_q = \mathbf{b}_q; \quad \boldsymbol{\rho}_q = \Lambda_q\boldsymbol{\rho},$$

74 where the multiple unknown vectors $\boldsymbol{\rho}_q$, $q = 1, \dots, S$, share the same support. The
 75 main advantage of this formulation is that we can immediately infer that the data
 76 vectors \mathbf{b}_q are linear combinations of the same M-columns of $\tilde{\mathcal{A}}$, those that belong to
 77 the support of the unknown $\boldsymbol{\rho}$. The implication is that the columns of $\tilde{\mathcal{A}}$ indexed by
 78 $T = \text{supp}(\boldsymbol{\rho})$ span the column subspace of B , the 'signal' subspace of B . Hence, the
 79 support T is the zero set of the orthogonal projections of the columns of matrix $\tilde{\mathcal{A}}$
 80 onto the null space of the data matrix B . Moreover, the support is recovered exactly
 81 under the assumption that all M-sets of columns of $\tilde{\mathcal{A}}$ are linearly independent. We
 82 discuss several imaging configurations for which the factorization (1.3) is feasible as
 83 well as instances where (1.3) holds only approximately and MUSIC is no longer exact
 84 even for noise free data.

Let us remark that for different excitations q we obtain multiple measurement
 vectors \mathbf{b}_q which correspond to linear transformations of the same unknown vector $\boldsymbol{\rho}$.
 The data can be arranged in a matrix $B \in \mathbb{C}^{N \times S}$ whose columns are the vectors \mathbf{b}_q ,
 and the MMV formulation may be expressed as a matrix-matrix equation

$$\tilde{\mathcal{A}}\mathbf{P} = B,$$

85 where the unknown is now the matrix $\mathbf{P} \in \mathbb{C}^{K \times S}$ whose columns are the vectors
 86 $\boldsymbol{\rho}_q = \Lambda_q\boldsymbol{\rho}$ that share the same support. The optimization can therefore be performed
 87 within the MMV formalism as described in [14, 24, 36, 37]. The main idea is to
 88 seek the solution with the minimal (2,1)-norm which consists in minimizing the ℓ_1
 89 norm of the vector formed by the ℓ_2 norms of the rows of the unknown matrix \mathbf{P} .
 90 This guarantees the common support of the solution's columns. We do not pursue
 91 this approach here and refer the reader to [12] for an application of this formalism
 92 to imaging strong scattering scenes as well as to [6] where an MMV formulation
 93 for synthetic aperture imaging of frequency and direction dependent reflectivity was
 94 introduced and analyzed.

95 We present several configurations in array imaging that can be cast under the
 96 general framework discussed here, such as single- and multiple-frequency array imag-
 97 ing using single- or multiple-receivers. All these problems can be formulated as (1.1)
 98 for a single measurement vector, or as (1.2) when multiple measurement vectors are
 99 available. We also consider the non-linear phase retrieval problem, which according
 100 to [31, 28, 29] can be reduced to a linear system of the form (1.2). This requires
 101 intensity data corresponding to multiple coherent illuminations which when using the
 102 polarization identity are transformed to interferometric data. We consider multiple
 103 frequency intensity data collected at a single receiver due to multiple coherent illumi-
 104 nations that could be generated by a spatial light modulator (SLM) [30]. The solution
 105 of (1.2) may then be computed with Single Receiver INTERferometry (SRINT) as in
 106 [29], ℓ_1 -minimization or MUSIC.

107 The performance of these imaging methods for the non-linear phase retrieval
 108 problem is studied with numerical simulations in an optical digital microscopy imaging
 109 regime. Our simulations allow us to assess the robustness of the different methods to
 110 modeling errors resulting to perturbations in the unknown phases of the recorded
 111 data. We consider phase perturbations that are either due to grid displacements or
 112 to wave propagation in a weakly inhomogeneous medium. Our conclusions are that
 113 SRINT provides the less satisfactory image in terms of resolution but it is the more

114 robust method when there are modeling errors, the ℓ_1 method has the best resolution
 115 but is not very robust with respect to noise, while MUSIC seems to be the more
 116 competitive method at moderate signal to noise ratio regimes because it has better
 117 resolution than SRINT and is less sensitive to noise than ℓ_1 -minimization.

118 The paper is organized as follows. In Section 2 we present in a abstract linear
 119 algebra framework the conditions under which ℓ_1 -minimization and MUSIC provide
 120 the exact solution to problems (1.1) and (1.2) respectively. We also analyze the
 121 performance of these methods for noisy data. In Section 3 we formulate the array
 122 imaging problem and consider some common configurations used in active array imag-
 123 ing. Moreover, we discuss how the imaging problem can be cast under the abstract
 124 framework of Section 2 and what are adequate data-structures to be used in imaging
 125 with ℓ_1 -minimization and MUSIC. In Section 4, we explore with numerical simulations
 126 the robustness of the imaging methods for the phase retrieval problem in an optical
 127 (digital) microscopy regime. In Section 5 we illustrate with numerical simulations how
 128 our abstract theoretical results are relevant in assessing image resolution. Section 6
 129 contains our conclusions.

130 **2. Linear algebra aspects of imaging algorithms.** In this section we dis-
 131 cuss under what conditions ℓ_1 -minimization and MUSIC algorithms provide the exact
 132 solution when there is no noise in the data. We also discuss the performance of these
 133 algorithms for noisy data. We assume that imaging can be formulated as a linear
 134 inverse problem of the form

$$135 \quad (2.1) \quad \mathcal{A}_l \boldsymbol{\rho} = \mathbf{b}_l,$$

136 that is underdetermined. In (2.1), the model matrix

$$137 \quad (2.2) \quad \mathcal{A}_l = \begin{pmatrix} \uparrow & \uparrow & & \uparrow \\ \mathbf{a}_1^{(l_1)} & \mathbf{a}_2^{(l_2)} & \dots & \mathbf{a}_K^{(l_K)} \\ \downarrow & \downarrow & & \downarrow \end{pmatrix} \in \mathbb{C}^{N \times K}$$

138 relates the unknown vector $\boldsymbol{\rho} \in \mathbb{C}^K$, which is the “image” to be constructed, to
 139 the transformed vector $\mathbf{b}_l \in \mathbb{C}^N$, which contains the data. This matrix is fixed by
 140 the physical setup of the imaging system and, therefore, it is given to us. However,
 141 the important observation here is that \mathcal{A}_l also depends on a parameter vector $\mathbf{l} =$
 142 $[l_1, l_2, \dots, l_K]^T$ which may be varied so as several transformed vectors \mathbf{b}_l of the same
 143 unknown $\boldsymbol{\rho}$ can be obtained.

144 If only one snapshot of array measurements is available for imaging, we solve
 145 (2.1) for a single measurement vector (SMV) \mathbf{l} using ℓ_1 minimization that promotes
 146 the assumed sparsity of the vector $\boldsymbol{\rho}$. In that case, we will write (2.1) simply as
 147 $\mathcal{A}\boldsymbol{\rho} = \mathbf{b}$. When several snapshots of array measurements corresponding to different
 148 parameter vectors \mathbf{l}_q are available, we solve the corresponding MMV problem using
 149 MUSIC. In that case, we will write (2.1) as $\mathcal{A}_q \boldsymbol{\rho} = \mathbf{b}_q$.

150 **2.1. ℓ_1 minimization-based methods.** In the imaging problems considered
 151 here we assume that the scatterers occupy only a small fraction of a region of interest
 152 called the image window IW. This means that the true reflectivity vector $\boldsymbol{\rho}_0$ is sparse,
 153 so the number of its entries that are different than zero, denoted by M , is much
 154 smaller than its length K . Thus, $M = |\text{supp}(\boldsymbol{\rho}_0)| \ll K$. This prior knowledge
 155 changes the imaging problem substantially because we can exploit the sparsity of $\boldsymbol{\rho}_0$
 156 by formulating (2.1) as an optimization problem which seeks the sparsest vector in

157 \mathbb{C}^K that equates model and data. Thus, for a single measurement vector \mathbf{b} we solve

158 (2.3)
$$\min \|\boldsymbol{\rho}\|_{\ell_1} \quad \text{subject to} \quad \mathcal{A}\boldsymbol{\rho} = \mathbf{b}.$$

159 In this form, we may be able to pick the true solution $\boldsymbol{\rho}_0$ if the matrix \mathcal{A} and the
 160 sparsity of $\boldsymbol{\rho}_0$ fulfill certain conditions. In particular, we have the following four
 161 theorems whose proofs are given in Appendix A. We denote by $\|\cdot\|_{\ell_2}$ and $\|\cdot\|_{\ell_1}$ the
 162 ℓ_2 and ℓ_1 norms of a vector, respectively.

163 THEOREM 2.1. *M-sparse solutions of $\mathcal{A}\boldsymbol{\rho} = \mathbf{b}$ are unique, if*

164 (2.4)
$$|\langle \mathbf{a}_i, \mathbf{a}_j \rangle| < \frac{1}{2M} \quad \forall i \neq j,$$

165 where we assume that the columns of matrix \mathcal{A} are normalized so that $\|\mathbf{a}_i\|_{\ell_2} = 1 \forall i$.

166 THEOREM 2.2. *The M-sparse solution of $\mathcal{A}\boldsymbol{\rho} = \mathbf{b}$ can be found as the solution of*

167 (2.5)
$$\min \|\boldsymbol{\eta}\|_{\ell_1}, \quad \text{subject to} \quad \mathcal{A}\boldsymbol{\eta} = \mathbf{b},$$

168 if

169 (2.6)
$$|\langle \mathbf{a}_i, \mathbf{a}_j \rangle| < \frac{1}{2M}, \quad \forall i \neq j,$$

170 where we assume that the columns of matrix \mathcal{A} are normalized so that $\|\mathbf{a}_i\|_{\ell_2} = 1 \forall i$.

171 THEOREM 2.3. *Let $\boldsymbol{\rho}$ be a solution of $\mathcal{A}\boldsymbol{\rho} = \mathbf{b}$, and let T be the index set of the
 172 support of $\boldsymbol{\rho}$, so*

173
$$T = \text{supp}(\boldsymbol{\rho}), \quad \text{and} \quad M = |T|.$$

174 Fix a positive $\varepsilon < 1/2$, and suppose that the matrix \mathcal{A} satisfies:

- 175 (i) The column vectors are normalized so that $\|\mathbf{a}_i\|_{\ell_2} = 1 \forall i$.
 176 (ii) The column vectors in the set T are approximately orthogonal, so

177 (2.7)
$$|\langle \mathbf{a}_i, \mathbf{a}_j \rangle| < \frac{\varepsilon}{M}, \quad \forall i, j \in T, i \neq j.$$

178 (iii) For any $j \in T$ the vicinity

179 (2.8)
$$S_j = \left\{ k \neq j \quad \text{s.t.} \quad |\langle \mathbf{a}_k, \mathbf{a}_j \rangle| \geq \frac{1}{2M} \right\}$$

180 has the properties

181 (2.9)
$$|\langle \mathbf{a}_k, \mathbf{a}_j \rangle| \leq 1 - 2\varepsilon \quad \forall k \in S_j,$$

182 and

183 (2.10)
$$|\langle \mathbf{a}_k, \mathbf{a}_j \rangle| < \frac{\varepsilon}{M} \quad \forall k \in S_i, \quad \forall i \neq j.$$

184 Then $\boldsymbol{\rho}$, the M-sparse solution of $\mathcal{A}\boldsymbol{\rho} = \mathbf{b}$, can be found as the solution of

185
$$\min \|\boldsymbol{\eta}\|_{\ell_1}, \quad \text{subject to} \quad \mathcal{A}\boldsymbol{\eta} = \mathbf{b}.$$

THEOREM 2.4. **Noisy case.** Let $\boldsymbol{\rho}$ be an M -sparse solution of

$$\mathcal{A}\boldsymbol{\rho} = \mathbf{b},$$

186 and let $T = \text{supp}(\boldsymbol{\rho})$, so $M = |T|$. Fix a positive $\varepsilon < 1/2$, and suppose that \mathcal{A} satisfies
 187 conditions (i), (ii), and (iii) of Theorem 2.3.

188 Furthermore, let $\boldsymbol{\rho}_\delta$ be the minimal ℓ_1 -norm solution of the noisy problem

$$189 \quad (2.11) \quad \min \|\boldsymbol{\eta}\|_{\ell_1}, \text{ subject to } \mathcal{A}\boldsymbol{\eta} = \mathbf{b}^\delta,$$

190 with \mathbf{b}^δ defined by

$$191 \quad (2.12) \quad \mathbf{b}^\delta = \mathbf{b} + \delta\mathbf{b},$$

192 such that the noise $\delta\mathbf{b}$ is bounded for some small positive δ , so that

$$193 \quad (2.13) \quad \|\delta\mathbf{b}\|_{\ell_2} \leq \delta.$$

194 Assume that \mathcal{A} has the property that the solution $\delta\boldsymbol{\rho}$ to

$$195 \quad (2.14) \quad \min \|\boldsymbol{\eta}\|_{\ell_1}, \text{ subject to } \mathcal{A}\boldsymbol{\eta} = \delta\mathbf{b},$$

196 satisfies

$$197 \quad (2.15) \quad \|\delta\boldsymbol{\rho}\|_{\ell_1} \leq C\|\delta\mathbf{b}\|_{\ell_2}.$$

198 Then, we can show that the solution $\boldsymbol{\rho}_\delta$ of (2.11) can be decomposed as

$$199 \quad (2.16) \quad \boldsymbol{\rho}_\delta = \boldsymbol{\rho}_c + \boldsymbol{\rho}_i,$$

200 with $\boldsymbol{\rho}_c$ the coherent part of the solution supported on T or in the vicinities S_j with
 201 $j \in T$, and $\boldsymbol{\rho}_i$ the incoherent part of the solution which is supported away from the
 202 vicinities and it is small. Specifically, for $\boldsymbol{\rho}_c$ we have that for any $j \in T$

$$203 \quad \left| |(\boldsymbol{\rho})_j| - |(\boldsymbol{\rho}_c)_j| + \sum_{k \in S_j} \langle \mathbf{a}_j, \mathbf{a}_k \rangle |(\boldsymbol{\rho}_c)_k| \right| \leq \delta_0 + C\delta,$$

204 with

$$205 \quad \delta_0 = \frac{2C\delta(1-\varepsilon)}{M(1-2\varepsilon)} + \frac{2\varepsilon(\|\boldsymbol{\rho}\|_{\ell_1} + C\delta)}{M}.$$

While for $\boldsymbol{\rho}_i$ we can show that:

$$\|\boldsymbol{\rho}_i\|_{\ell_1} \leq \delta_1,$$

206 with δ_1 given by

$$207 \quad \delta_1 = C\delta + \frac{4C\delta(1-\varepsilon)}{(1-2\varepsilon)}.$$

208 Theorems 2.1 and 2.2 are well known results in the literature of compressive
 209 sensing [21, 16, 22]. The first theorem tells us that the M -sparse solution of the linear
 210 system $\mathcal{A}\boldsymbol{\rho} = \mathbf{b}$ is unique when the columns of the matrix satisfy the orthonormality
 211 condition (2.4). This condition is satisfied when the mutual coherence of the matrix
 212 \mathcal{A} , defined as $\max_{i \neq j} |\langle \mathbf{a}_i, \mathbf{a}_j \rangle|$, is smaller than $1/(2M)$. This first theorem is an
 213 ℓ_0 uniqueness result. The second result, Theorem 2.2, tells us that the unique M -
 214 sparse solution of $\mathcal{A}\boldsymbol{\rho} = \mathbf{b}$ can be found by solving the ℓ_1 minimization problem (2.5).

Algorithm 1 GelMa for solving (2.5)

Require: Set $\mathbf{y} = \mathbf{0}$, $\mathbf{z} = \mathbf{0}$. Pick the step size β , and a regularization parameter τ .

repeat

 Compute the residual $\mathbf{r} = \mathbf{b} - \mathcal{A}\mathbf{y}$

$\mathbf{y} \leftarrow \eta_{\tau\beta}(\mathbf{y} + \beta\mathcal{A}^*(\mathbf{z} + \mathbf{r}))$

$\mathbf{z} \leftarrow \mathbf{z} + \beta\mathbf{r}$

until Convergence

215 This is a very useful result because it is the ℓ_1 minimization problem that can be
216 solved efficiently in practice, for example, by using the algorithm GelMa described in
217 Algorithm 1, which involves only simple matrix-vector multiplications followed by a
218 shrinkage-thresholding step defined by the operator $\eta_\tau(y_i) = \text{sign}(y_i) \max\{0, |y_i| - \tau\}$.
219 In the noiseless case, this algorithm converges to the exact solution independently of
220 the value of the regularization parameter τ . For more details we refer to [27].

221 Theorem 2.3 is to the best of our knowledge new. Its proof is given in Appendix
222 A. This theorem tells us that the M-sparse solution of $\mathcal{A}\boldsymbol{\rho} = \mathbf{b}$ can be recovered
223 by solving the ℓ_1 minimization problem under a less stringent condition than (2.6)
224 provided that the column vectors of the matrix \mathcal{A} that are in the support of the true
225 solution $\boldsymbol{\rho}_0$ are approximately orthogonal, that is, they satisfy (2.7). Note that we
226 allow for the columns of \mathcal{A} to be close to collinear. Moreover, we define the vicinities
227 S_j for the column vectors \mathbf{a}_j in the support of the true solution, and we assume that
228 all the column vectors that are in the vicinity of a *support column vector* are close
229 enough to it, so (2.9) holds. We also assume that the vicinities S_i and S_j , for $i \neq j$,
230 are far enough, so (2.10) holds.

231 The last result, Theorem 2.4, is the noisy version of Theorem 2.3. It shows that
232 when the data \mathbf{b} is not exact but is known up to some bounded vector $\boldsymbol{\delta}\mathbf{b}$, the solution
233 $\boldsymbol{\rho}_\delta$ of the minimization problem (2.11)-(2.12) is close to the solution of the original
234 (noiseless) problem in the following sense. The solution $\boldsymbol{\rho}_\delta$ can be decomposed in
235 two parts: the coherent part $\boldsymbol{\rho}_c$ supported in T or in the vicinities S_j , $j \in T$, of the
236 true solution, and the incoherent part $\boldsymbol{\rho}_i$ usually referred to as grass in imaging. The
237 grass is supported away from the vicinities S_j and it is shown to be small assuming
238 that (2.15) holds for the solution to (2.14) and assuming that the norm of the noise
239 is small so (2.13) holds. Other stability results can be found in [7, 8, 17, 35, 18, 4].

240 We will see in Section 5 how Theorems 2.3 and 2.4 can be applied in imaging.

241 **2.2. MUSIC.** MUSIC is a subspace imaging algorithm based on the decomposi-
242 tion of the measurements into two orthogonal domains: the signal and noise subspaces
243 [34]. The key is to be able to form a data matrix

$$244 \quad (2.17) \quad B = \begin{pmatrix} b_{11} & b_{12} & \dots & b_{1S} \\ b_{21} & b_{22} & \dots & b_{2S} \\ \dots & \dots & \dots & \dots \\ b_{N1} & b_{N2} & \dots & b_{NS} \end{pmatrix} = \begin{pmatrix} \uparrow & \uparrow & \dots & \uparrow \\ \mathbf{b}_1 & \mathbf{b}_2 & \dots & \mathbf{b}_S \\ \downarrow & \downarrow & \dots & \downarrow \end{pmatrix} \in \mathbb{C}^{N \times S},$$

245 whose column vectors \mathbf{b}_q are obtained from a family of linear systems $\mathcal{A}_q\boldsymbol{\rho} = \mathbf{b}_q$ that
246 can be rewritten in the form

$$247 \quad (2.18) \quad \tilde{\mathcal{A}}\Lambda_q\boldsymbol{\rho} = \mathbf{b}_q, \quad q = 1, \dots, S,$$

248 where Λ_q is a diagonal matrix whose entries can be controlled to form the images.
249 The assumption here is that the model matrices \mathcal{A}_q relating the unknown vector $\boldsymbol{\rho}$

250 with the data vectors \mathbf{b}_q can be factorized into two matrices
 (2.19)

$$251 \quad \tilde{\mathcal{A}} = \begin{pmatrix} \uparrow & \uparrow & & \uparrow \\ \tilde{\mathbf{a}}_1 & \tilde{\mathbf{a}}_2 & \dots & \tilde{\mathbf{a}}_K \\ \downarrow & \downarrow & & \downarrow \end{pmatrix} \in \mathbb{C}^{N \times K} \quad \text{and} \quad \Lambda_q = \begin{pmatrix} l_{1q} & 0 & & \\ 0 & l_{2q} & & \\ & & \ddots & \\ & & & 0 & l_{Kq} \end{pmatrix} \in \mathbb{C}^{K \times K},$$

252 with $\tilde{\mathcal{A}}$ independent of the parameter vector $\mathbf{l}_q = [l_{1q}, l_{2q}, \dots, l_{Kq}]^T$, and Λ_q
 253 diagonal. Under this assumption, the imaging problem (2.18) can be reinterpreted in
 254 the form of an MMV problem

$$255 \quad (2.20) \quad \tilde{\mathcal{A}}\boldsymbol{\rho}_q = \mathbf{b}_q,$$

256 with $\boldsymbol{\rho}_q = \Lambda_q\boldsymbol{\rho}$. Physically, each $\boldsymbol{\rho}_q$ is a transformed version of the same unknown
 257 vector $\boldsymbol{\rho}$. The data can be arranged into the data matrix (2.17), and (2.20) may be
 258 expressed as a matrix-matrix equation

$$259 \quad (2.21) \quad \tilde{\mathcal{A}}\mathbf{P} = B,$$

260 where the columns of $\mathbf{P} \in \mathbb{C}^{K \times S}$, $\boldsymbol{\rho}_q = \Lambda_q\boldsymbol{\rho}$, share the same support.

261 The important element of the new formulation (2.20) (or (2.21)) is that now all
 262 the data vectors \mathbf{b}_q are linear combinations of the same M columns of $\tilde{\mathcal{A}}$ (or \mathcal{A}), those
 263 columns that correspond to $T = \text{supp}(\boldsymbol{\rho})$, with $M = |T|$. Thus, every column of $\tilde{\mathcal{A}}$
 264 indexed by T is contained in the column space of B , the signal subspace, which is
 265 orthogonal to the noise subspace. Hence, one can simply find the unknown support
 266 T by projecting the columns of $\tilde{\mathcal{A}}$ onto the noise subspace. Both, the signal and the
 267 noise subspaces can be obtained via the singular value decomposition (SVD) of B .

268 More precisely, the objective of a MUSIC algorithm is to find the support T of
 269 an unknown sparse vector $\boldsymbol{\rho} = [\rho_1, \rho_2, \dots, \rho_K]^T$ with a number of nonzero entries M
 270 much smaller than its length K . With a sufficiently diverse number of experiments
 271 $S \geq M$ we create a data matrix B , and we compute its SVD

$$272 \quad (2.22) \quad B = U\Sigma V^* = \sum_{j=1}^K \sigma_j \mathbf{u}_j \mathbf{v}_j^*.$$

273 If the data is noiseless there are M nonzero singular values $\sigma_1 > \sigma_2 > \dots > \sigma_M > 0$
 274 with corresponding (left) singular vectors \mathbf{u}_j , $j = 1, \dots, M$ that span the signal
 275 subspace of \mathbb{C}^N . The remaining singular values σ_j , $j = M + 1, \dots, K$, are zero, and
 276 the corresponding (left) singular vectors span the noise subspace of \mathbb{C}^N . Because the
 277 set of columns of $\tilde{\mathcal{A}}$ indexed by $T = \text{supp}(\boldsymbol{\rho})$ also spans the signal subspace, the sought
 278 support T corresponds to the zero set of the orthogonal projections of the columns
 279 vectors $\tilde{\mathbf{a}}_k$ onto the noise subspace. Thus, it follows that the support of $\boldsymbol{\rho}$ can be
 280 found among the zeros of the imaging functional

$$281 \quad (2.23) \quad \mathcal{I}_k^{\text{SIGNAL}} = \sum_{j=1}^M |\tilde{\mathbf{a}}_k^* \mathbf{u}_j|^2, \quad k = 1, \dots, K,$$

282 or, equivalently, among the peaks of the imaging functional

$$283 \quad (2.24) \quad \mathcal{I}_k^{\text{MUSIC}} = \frac{\|\tilde{\mathbf{a}}_k\|_{\ell_2}}{\sum_{j=M+1}^N |\tilde{\mathbf{a}}_k^* \mathbf{u}_j|^2}, \quad k = 1, \dots, K.$$

284 Furthermore, if all sets of M columns of \tilde{A} are linearly independent, then the peaks
 285 exactly coincide with the support of $\boldsymbol{\rho}$ in the noiseless case. In (2.24), the numerator
 286 is a normalization factor.

287 Once the support is recovered, the problem typically becomes overdetermined
 288 ($N > |\text{supp}(\boldsymbol{\rho})|$) and the nonzero values of $\boldsymbol{\rho}$ can be easily found by solving the linear
 289 system restricted to the given support with an ℓ_2 method [13].

290 Regarding imaging with noisy data, it follows from Weyl's theorem [39] that when
 291 noise is added to the data so $B \rightarrow B^\delta = B + E$ with $\|E\|_{\ell_2} < \delta$, then no singular
 292 value σ^δ moves more than the norm of the perturbation, i.e., $\|\sigma^\delta - \sigma\|_{\ell_2} < \delta$. Hence,
 293 (i) perturbed and unperturbed singular values are paired, and (ii) the spectral gap
 294 between the zero and the nonzero singular values remains large if the smallest nonzero
 295 unperturbed singular value $\sigma_M \gg \delta$. If the noise is not too large, then the rank of
 296 the data matrix B^δ can be determined, and so is $M = |T|$.

297 The signal and noise subspaces are also perturbed in the presence of noise. It can
 298 be shown, however, that the perturbed subspaces remain close to the unperturbed
 299 ones, with changes that are proportional to the reciprocal of the spectral gap $\beta =$
 300 $\sigma_M^\delta - \sigma_{M+1}$. This follows from Wedin's Theorem [38].

301 **THEOREM 2.5.** (Wedin) *Let B have the SVD $B = Q + Q_0$ with $Q = U\Sigma V^T$ and*
 302 *$Q_0 = U_0\Sigma_0V_0^T$, and let the perturbed matrix $B^\delta = B + E$ have the SVD $B^\delta = Q^\delta + Q_0^\delta$*
 303 *with $Q^\delta = U^\delta\Sigma^\delta V^{\delta T}$ and $Q_0^\delta = U_0^\delta\Sigma_0^\delta V_0^{\delta T}$. If there exist two constants $\alpha \geq 0$ and*
 304 *$\beta > 0$ such that $\sigma_{\max}(Q_0) \leq \alpha$ and $\sigma_{\min}(Q^\delta) \geq \alpha + \beta$, then the distance between the*
 305 *orthogonal projections onto the subspaces $R(Q)$ and $R(Q^\delta)$ is bounded by*

$$306 \quad (2.25) \quad \|P_{R(Q^\delta)} - P_{R(Q)}\|_{\ell_2} \leq \frac{\delta}{\beta},$$

307 where $\delta = \max(\|EV\|_{\ell_2}, \|E^*U\|_{\ell_2})$.

308 There is much work done on the robustness of MUSIC with respect to noise. We
 309 refer to [26], and references therein, for a recent discussion about how much noise
 310 the MUSIC algorithm can tolerate. When we apply the Theorem 2.5 to our imaging
 311 problem, where $Q_0 = 0$, we obtain the following result whose proof is in Appendix B.
 312

313 **THEOREM 2.6.** *Let $X = \text{Diag}(\boldsymbol{\rho})$ be a diagonal matrix that solves*

$$314 \quad (2.26) \quad \tilde{A}XL = B,$$

315 where \tilde{A} satisfies conditions (i), (ii), and (iii) of Theorem 2.3 for a fixed $\varepsilon < 1/3$,

$$316 \quad L = \begin{pmatrix} l_{11} & l_{12} & l_{1S} \\ l_{21} & l_{22} & l_{2S} \\ \vdots & \vdots & \vdots \\ l_{K1} & l_{K2} & l_{KS} \end{pmatrix} \in \mathbb{C}^{K \times S},$$

317 and B is the noiseless data matrix (2.17) with SVD $B = Q = U\Sigma V^T$. Let the
 318 perturbed matrix $B^\delta = Q^\delta + Q_0$ be such that $\sigma_{\max}(B^\delta - B) \leq \delta$. Suppose $\boldsymbol{\rho}$, the vector
 319 diagonal entries of X , is sparse with $T = \text{supp}(\boldsymbol{\rho})$, $M = |T|$, $M \ll \text{size}(\boldsymbol{\rho})$, and

$$320 \quad \rho_m = \min_{\rho_i \neq 0} \{|\rho_i|\}.$$

321 Let L_T be the submatrix of L , formed by the rows corresponding to T , has

$$322 \quad (2.27) \quad \sigma_m^T = \sigma_{\min}(L_T).$$

323 If

$$324 \quad (2.28) \quad 2\delta < \rho_m \sigma_m^T (1 - 3\varepsilon),$$

325 the orthogonal projections onto the subspaces $R(Q^\delta)$ and $R(B)$ are close:

$$326 \quad (2.29) \quad \|P_{R(Q^\delta)} - P_{R(B)}\|_{\ell_2} \leq \frac{\delta}{\rho_m \sigma_m^T (1 - 3\varepsilon)}.$$

327 To conclude, the main step in setting up MUSIC is to be able to find a suitable
 328 factorization of the model matrix as $\mathcal{A}_q = \tilde{\mathcal{A}}\Lambda_q$, where Λ_q is diagonal. In that case,
 329 the imaging vectors are just the columns of $\tilde{\mathcal{A}}$ that are given. We discuss next imaging
 330 situations in which this factorization is possible and MUSIC can form images
 331 with high precision. We also discuss applications in which the factorization is only
 332 approximate and, hence, images obtained with MUSIC lose resolution.

333 **3. Array imaging: data models.** The goal of array imaging is to form images
 334 inside a region of interest called the image window IW. In active array imaging the
 335 array probes the medium by sending signals and recording the echoes. Probing of the
 336 medium can be done with many different types of arrays that differ in their number
 337 of transmitters and receivers, their geometric layouts, or the type of signals they use
 338 for illumination. They may use single frequency signals sent from different positions,
 339 or multifrequency signals sent from one or more positions. Of course, the problem of
 340 active array imaging also depends on the receivers. They can record the intensities
 341 and phases of the signals that arrive to the array or only their intensities.

342 In this section, we describe some common configurations used in active array
 343 imaging. The array, with N transducers separated by a distance h , has a characteristic
 344 length a (see Figure 1). The transducers emit signals from positions \vec{x}_s and record
 345 the echoes at positions \vec{x}_r , $s, r = 1, 2, \dots, N$. They can use single or multifrequency
 346 signals, with frequencies ω_l , $l = 1, \dots, S$. Our goal is to reconstruct a sparse scene
 347 consisting of M point-scatterers at a distance L from the array, whose positions \vec{y}_{n_j}
 348 and reflectivities $\alpha_{n_j} \in \mathbb{C}$, $j = 1, \dots, M$, we seek to determine. The ambient medium
 349 between the array and the scatterers can be homogeneous or inhomogeneous.

In order to form the images we discretize the IW using a uniform grid of points
 \vec{y}_k , $k = 1, \dots, K$, and we introduce the *true reflectivity vector*

$$\boldsymbol{\rho}_0 = [\rho_{01}, \dots, \rho_{0K}]^T \in \mathbb{C}^K,$$

350 such that $\rho_{0k} = \sum_{j=1}^M \alpha_{n_j} \delta_{\vec{y}_{n_j} \vec{y}_k}$, $k = 1, \dots, K$, where δ_{\cdot} is the classical Kronecker
 351 delta. We will not assume that the scatterers lie on the grid, i.e., $\{\vec{y}_{n_1}, \dots, \vec{y}_{n_M}\} \not\subset$
 352 $\{\vec{y}_1, \dots, \vec{y}_K\}$ in general. To write the data received on the array in a compact form,
 353 we define the Green's function vector

$$354 \quad (3.1) \quad \hat{\mathbf{g}}(\vec{y}; \omega) = [\hat{G}(\vec{x}_1, \vec{y}; \omega), \hat{G}(\vec{x}_2, \vec{y}; \omega), \dots, \hat{G}(\vec{x}_N, \vec{y}; \omega)]^T$$

355 at location \vec{y} in the IW, where $\hat{G}(\vec{x}, \vec{y}; \omega)$ denotes the free-space Green's function of the
 356 homogeneous or inhomogeneous medium. This function characterizes the propagation
 357 of a signal of angular frequency ω from point \vec{y} to point \vec{x} , so (3.1) represents the
 358 signal received at the array due to a point source of frequency ω at \vec{y} . When the
 359 medium is homogeneous,

$$360 \quad (3.2) \quad \hat{G}(\vec{x}, \vec{y}; \omega) = \hat{G}_0(\vec{x}, \vec{y}; \omega) = \frac{\exp(i\kappa|\vec{x} - \vec{y}|)}{4\pi|\vec{x} - \vec{y}|}, \quad \kappa = \frac{\omega}{c_0}.$$

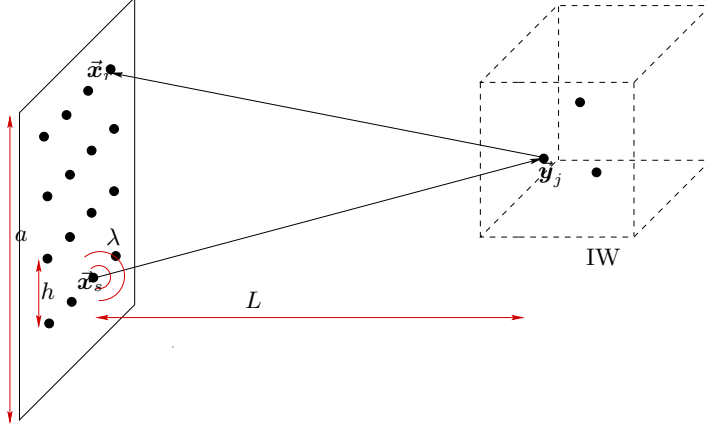


FIG. 1. General setup of an array imaging problem. The transducer at \vec{x}_s emits a probing signal and the reflected intensities are recorded at \vec{x}_r . The scatterers located at \vec{y}_j , $j = 1, \dots, M$ are at distance L from the array and inside the image window IW .

361 In this case, the Green's function vector is

$$362 \quad \hat{\mathbf{g}}_0(\vec{y}; \omega) = [\hat{G}_0(\vec{x}_1, \vec{y}; \omega), \hat{G}_0(\vec{x}_2, \vec{y}; \omega), \dots, \hat{G}_0(\vec{x}_N, \vec{y}; \omega)]^T.$$

363 We assume that the scatterers are far apart or that the reflectivities are small, so
 364 multiple scattering between them is negligible. In this case, the Born approximation
 365 holds and, thus, the response at \vec{x}_r (including phases) due to a pulse of angular
 366 frequency ω_l sent from \vec{x}_s , and reflected by the M scatterers, is given by

$$367 \quad (3.3) \quad P(\vec{x}_r, \vec{x}_s; \omega_l) = \sum_{j=1}^M \alpha_j G(\vec{x}_r, \vec{y}_{n_j}; \omega_l) G(\vec{y}_{n_j}, \vec{x}_s; \omega_l),$$

368 and the the full response matrix that contains all possible information for imaging by

$$369 \quad (3.4) \quad P(\omega_l) = [P(\vec{x}_r, \vec{x}_s; \omega_l)] = \sum_{j=1}^M \alpha_j \hat{\mathbf{g}}(\vec{y}_{n_j}; \omega_l) \hat{\mathbf{g}}^T(\vec{y}_{n_j}; \omega_l).$$

370 Next, we describe different situations of interest in active array imaging.

371 **3.1. Single frequency signals and multiple receivers.** Let us first consider
 372 the case in which only one illumination of frequency ω is sent using the N sources in
 373 the array located at positions \vec{x}_s , $s = 1, \dots, N$. The echoes are also recorded at the
 374 N receivers located at \vec{x}_r , $r = 1, \dots, N$. If $\hat{\mathbf{f}}(\omega) = [\hat{f}_1(\omega), \dots, \hat{f}_N(\omega)]^T$ represents the
 375 illumination vector whose entries are the signals sent from the sources in the array,
 376 then $\hat{g}_{\hat{\mathbf{f}}(\omega)}^{(k)} = \hat{\mathbf{g}}(\vec{y}_k; \omega)^T \hat{\mathbf{f}}(\omega)$ is the field at the grid position \vec{y}_k in the IW. Thus,

$$377 \quad (3.5) \quad \mathcal{A}_{\hat{\mathbf{f}}(\omega)} = \begin{pmatrix} \uparrow & & & \uparrow & & & \uparrow \\ \hat{g}_{\hat{\mathbf{f}}(\omega)}^{(1)} \hat{\mathbf{g}}(\vec{y}_1; \omega) & \hat{g}_{\hat{\mathbf{f}}(\omega)}^{(2)} \hat{\mathbf{g}}(\vec{y}_2; \omega) & \dots & \hat{g}_{\hat{\mathbf{f}}(\omega)}^{(K)} \hat{\mathbf{g}}(\vec{y}_K; \omega) & & & \\ \downarrow & & & \downarrow & & & \downarrow \end{pmatrix} \in \mathbb{C}^{N \times K}$$

378 is the model matrix that connects the unknown reflectivity vector $\boldsymbol{\rho} \in \mathbb{C}^K$ to the data
 379 vector $\mathbf{b}_{\hat{\mathbf{f}}(\omega)} \in \mathbb{C}^N$ that depends on the illumination $\hat{\mathbf{f}}(\omega)$.

380 If a single illumination is used to form an image, then active array imaging
 381 amounts to finding $\boldsymbol{\rho}$ from the system of linear equations

$$382 \quad (3.6) \quad \mathcal{A}_{\widehat{\mathbf{f}}(\omega)} \boldsymbol{\rho} = \mathbf{b}_{\widehat{\mathbf{f}}(\omega)}.$$

383 Abusing a little bit the notation used in Section 2, we have indicated in (3.6) that
 384 the control parameter vector is the illumination $\widehat{\mathbf{f}}(\omega)$. According to (2.1)-(2.2), the
 385 parameter vector is $\mathbf{l} = [\widehat{\mathbf{g}}_{\widehat{\mathbf{f}}(\omega)}^{(1)}, \widehat{\mathbf{g}}_{\widehat{\mathbf{f}}(\omega)}^{(2)}, \dots, \widehat{\mathbf{g}}_{\widehat{\mathbf{f}}(\omega)}^{(K)}]^T$ which depends on the Green's func-
 386 tion vectors $\widehat{\mathbf{g}}(\widehat{\mathbf{y}}; \omega)$ fixed by the physical layout, and on the illumination vector $\widehat{\mathbf{f}}(\omega)$
 387 that we control. The system of linear equations (3.6) can be solved using appropriate
 388 ℓ_2 or ℓ_1 methods. If an ℓ_1 -norm minimization method is chosen, we would seek the
 389 sparsest vector $\boldsymbol{\rho}$ among all possible vectors satisfying (3.6).

390 If, instead, multiple illuminations are used to form the images, then we can use
 391 an MMV approach to find the solution with MUSIC. Indeed, note that the model
 392 matrix (3.5) can be factorized into two matrices

$$393 \quad (3.7) \quad \tilde{\mathcal{A}} = \begin{pmatrix} \uparrow & \uparrow & & \uparrow \\ \widehat{\mathbf{g}}(\widehat{\mathbf{y}}_1; \omega) & \widehat{\mathbf{g}}(\widehat{\mathbf{y}}_2; \omega) & \dots & \widehat{\mathbf{g}}(\widehat{\mathbf{y}}_K; \omega) \\ \downarrow & \downarrow & & \downarrow \end{pmatrix} \in \mathbb{C}^{N \times K}$$

394 and

$$395 \quad (3.8) \quad \Lambda_{\widehat{\mathbf{f}}(\omega)} = \begin{pmatrix} \widehat{\mathbf{g}}_{\widehat{\mathbf{f}}(\omega)}^{(1)} & 0 & & \\ 0 & \widehat{\mathbf{g}}_{\widehat{\mathbf{f}}(\omega)}^{(2)} & & \\ & & \ddots & \\ & & 0 & \widehat{\mathbf{g}}_{\widehat{\mathbf{f}}(\omega)}^{(k)} \end{pmatrix} \in \mathbb{C}^{K \times K},$$

396 so that $\mathcal{A}_{\widehat{\mathbf{f}}(\omega)} = \tilde{\mathcal{A}} \Lambda_{\widehat{\mathbf{f}}(\omega)}$. Hence, it follows from the discussion in Section 2 that (3.6)
 397 can be written in the MMV form

$$398 \quad (3.9) \quad \tilde{\mathcal{A}} \tilde{\boldsymbol{\rho}}_q = \mathbf{b}_q, \quad q = 1, \dots, S,$$

399 and the support of $\boldsymbol{\rho}$ can be found exactly with MUSIC if enough data vectors $\mathbf{b}_{\widehat{\mathbf{f}}_q(\omega)}$
 400 are available. In (3.9), $\mathbf{b}_q = \mathbf{b}_{\widehat{\mathbf{f}}_q(\omega)}$, and $\tilde{\boldsymbol{\rho}}_q = \Lambda_{\widehat{\mathbf{f}}_q(\omega)} \boldsymbol{\rho}$ represents an *effective source*
 401 *weighted reflectivity vector* with the same support as $\boldsymbol{\rho}$, and whose nonzero entries
 402 vary with $\widehat{\mathbf{f}}_q(\omega)$. We remark that the equivalent source problem (3.9) can be used to
 403 account for multiple scattering between the scatterers (see [12] for details).

404 To show that Theorem 2.6 is relevant for imaging we write (3.9) as (2.26) with
 405 the unknown matrix $X = \text{Diag}(\boldsymbol{\rho})$, the data matrix B formed by the S vectors \mathbf{b}_q , and
 406 the illumination matrix

$$407 \quad L = \begin{pmatrix} \uparrow & \uparrow & & \uparrow \\ \tilde{\mathcal{A}}^T \widehat{\mathbf{f}}_1(\omega) & \tilde{\mathcal{A}}^T \widehat{\mathbf{f}}_2(\omega) & \dots & \tilde{\mathcal{A}}^T \widehat{\mathbf{f}}_S(\omega) \\ \downarrow & \downarrow & & \downarrow \end{pmatrix} \in \mathbb{C}^{K \times S}$$

408 whose i th column $\tilde{\mathcal{A}}^T \widehat{\mathbf{f}}_i(\omega) = [\widehat{\mathbf{g}}_{\widehat{\mathbf{f}}_i(\omega)}^{(1)}, \widehat{\mathbf{g}}_{\widehat{\mathbf{f}}_i(\omega)}^{(2)}, \dots, \widehat{\mathbf{g}}_{\widehat{\mathbf{f}}_i(\omega)}^{(K)}]^T$ contains the fields at all grid
 409 positions $\widehat{\mathbf{y}}_k$, $k = 1, \dots, K$ due to the illumination $\widehat{\mathbf{f}}_i(\omega)$. Then, condition (2.27) can
 410 be interpreted as an orthogonality condition on the illuminations. Furthermore, if we

411 suppose that $S = N$ and use the illuminations $\hat{\mathbf{f}}_q(\omega) = \hat{f}(\omega)\hat{\mathbf{e}}_q$ ($\hat{\mathbf{e}}_q$ is the vector with
412 a 1 in the q th coordinate and 0's elsewhere) for all $q = 1, \dots, S$, then $L = \hat{f}(\omega)\tilde{\mathcal{A}}^T$.
413 In this case, $\sigma_m^T = \sigma_{\min}(L_T) \geq (1 - 3\varepsilon)|\hat{f}(\omega)|$, assuming $\tilde{\mathcal{A}}$ satisfies conditions (i), (ii)
414 and (iii) of Theorem 2.3 (see proof of Theorem 2.6 in Appendix B).

415 3.2. Multifrequency signals and one receiver: the one-dimensional prob-

416 **lem.** Consider now a one-dimensional problem with scatterers located at different
417 ranges. To determine their positions we only use one transducer that emits and re-
418 ceives multiple frequency signals. We assume that the scatterers are far from the
419 transducer, but not far from each other so the denominator of the Green's function
420 in (3.2) can be approximated by a constant. In that case, the collected data are
421 approximately the Fourier transform of the reflectivity vector to be imaged.

422 To fix ideas, denote by $z_n = L + (n - 1)\Delta z$ the distance between the single
423 transducer and the scatterer of reflectivity ρ_n , $n = 1, \dots, K$. Then,

$$424 \quad (3.10) \quad \sum_{n=1}^K e^{i2\kappa_m z_n} \rho_n = b_m, \quad m = 1, \dots, 2S,$$

425 relates the positions and reflectivities of the scatterers to the measurements b_m at
426 frequencies $\omega_m = \kappa_m c_0$, where c_0 is the wave speed in a homogeneous medium. In
427 this problem, we seek to recover the unknown vector $\boldsymbol{\rho} = [\rho_1, \rho_2, \dots, \rho_K]$ from the
428 multifrequency data vector $\mathbf{b} = [b_1, b_2, \dots, b_{2S}]$ recorded at a single receiver.

429 The next assumption allows to succinctly formulate one-dimensional multifre-
430 quency MUSIC in the form of an MMV problem using the Prony-type argument (see,
431 for example, [25]). Namely, suppose that the measurements are obtained at equally
432 spaced (spatial) frequencies $\kappa_m = \kappa_1 + (m - 1)\Delta\kappa$, $m = 1, 2, \dots, 2S$. Then, we write
433 (3.10) in matrix form as

$$434 \quad (3.11) \quad \mathcal{A}_{2S} \boldsymbol{\rho} = \mathbf{b},$$

435 where

$$436 \quad (3.12) \quad \mathcal{A}_{2S} = \begin{pmatrix} e^{i2\kappa_1 z_1} & e^{i2\kappa_1 z_2} & \dots & e^{i2\kappa_1 z_K} \\ e^{i2\kappa_2 z_1} & e^{i2\kappa_2 z_2} & \dots & e^{i2\kappa_2 z_K} \\ \dots & \dots & \dots & \dots \\ e^{i2\kappa_{2S} z_1} & e^{i2\kappa_{2S} z_2} & \dots & e^{i2\kappa_{2S} z_K} \end{pmatrix}$$

437 is a Vandermonde matrix of dimensions $2S \times K$. Since we only have one data vector
438 $\mathbf{b} \in \mathbb{C}^{2S}$ we cannot determine from it a signal space of dimension $M = |\text{supp}(\boldsymbol{\rho})|$.
439 However, following the general idea of Prony-type [32] methods we form the $S \times S$
440 data matrix

$$441 \quad (3.13) \quad B = \begin{pmatrix} b_1 & b_2 & \dots & b_S \\ b_2 & b_3 & \dots & b_{S+1} \\ \dots & \dots & \dots & \dots \\ b_S & b_{S+1} & \dots & b_{2S} \end{pmatrix},$$

442 whose rank is M if $S > M$. If we now set the $S \times K$ matrix

$$443 \quad (3.14) \quad \tilde{\mathcal{A}} = \mathcal{A}_S = \begin{pmatrix} e^{i2\kappa_1 z_1} & e^{i2\kappa_1 z_2} & \dots & e^{i2\kappa_1 z_K} \\ e^{i2\kappa_2 z_1} & e^{i2\kappa_2 z_2} & \dots & e^{i2\kappa_2 z_K} \\ \dots & \dots & \dots & \dots \\ e^{i2\kappa_S z_1} & e^{i2\kappa_S z_2} & \dots & e^{i2\kappa_S z_K} \end{pmatrix}$$

444 and the $K \times K$ diagonal matrices

$$445 \quad (3.15) \quad \Lambda_q = \begin{pmatrix} e^{i2\Delta\kappa z_1} & 0 & \dots & 0 & 0 \\ 0 & e^{i2\Delta\kappa z_2} & \dots & 0 & 0 \\ \dots & \dots & \dots & e^{i2\Delta\kappa z_{K-1}} & 0 \\ 0 & 0 & \dots & 0 & e^{i2\Delta\kappa z_K} \end{pmatrix}^q,$$

with $q = 1, \dots, S$, then it is straightforward to verify that $\tilde{\mathcal{A}}\Lambda_q\boldsymbol{\rho} = \mathbf{b}_q$, where \mathbf{b}_q is the q th column of the matrix B in (3.13). Thus, we obtain the desired structure

$$\tilde{\mathcal{A}}\boldsymbol{\rho}_q = \mathbf{b}_q,$$

446 and MUSIC can be applied directly to find the support of $\boldsymbol{\rho}$. Subsequently, as noted
447 above $\boldsymbol{\rho}$ itself can be determined by solving the linear system restricted on the support
448 $\boldsymbol{\rho}$.

449 If $M \ll K$, so the vector $\boldsymbol{\rho}$ is M -sparse, then the solution can also be found
450 directly from (3.11) by using an ℓ_1 -norm minimization approach. Note that (3.11)
451 always has a unique M -sparse solution if $M < S$. Indeed, we argue by contradiction
452 that it is not possible to have more than one M -sparse solution if $M < S$. Suppose
453 there are two M -sparse solutions $\boldsymbol{\rho}_1$ and $\boldsymbol{\rho}_2$. Then, $\mathcal{A}_{2S}\mathbf{y} = 0$ for $\mathbf{y} = \boldsymbol{\rho}_1 - \boldsymbol{\rho}_2$. Since
454 the support of \mathbf{y} is less or equal than $2M$, we have $2M$ linearly dependent columns
455 of \mathcal{A}_{2S} , which is impossible for Vandermonde matrices since they are full rank.

456 **3.3. The single frequency phase retrieval problem.** In its classical form,
457 the phase retrieval problem consists in finding a function h from the amplitude of its
458 Fourier transform \hat{h} . In imaging, it consists in finding a vector $\boldsymbol{\rho}$ that is compatible
459 with a set of quadratic equations for measured amplitudes. This occurs in imaging
460 regimes where only intensity data is recorded, which means that most of the infor-
461 mation encoded in the phases is lost. Phase retrieval algorithms have been developed
462 over a long time to deal with this problem [20, 19]. They are flexible and effective
463 but depend on prior information about the image and can give uneven results. An
464 alternative convex approach that guarantees exact recovery has been considered in
465 [10, 9] but its computational cost is extremely high when the problem is large. When,
466 however, multiple measurements of the object to be imaged are available, we may re-
467 cover the missing phase information and image holographically much more efficiently
468 [31, 28, 29]. By holographic imaging we mean the use of interference patterns between
469 two or more coherent sources in order to form the images [40].

Indeed, let us consider single frequency imaging with multiple sources and re-
ceivers as in problem (3.9), where the data vectors $\mathbf{b}_q = \tilde{\mathcal{A}}\tilde{\boldsymbol{\rho}}_q$, that depend on the
illumination $\hat{\mathbf{f}}_q(\omega)$, contained the amplitudes and phases of the recorded signals We
now, however, assume that only the amplitudes squared of the components of these
data vectors can be measured. Then, the phase retrieval problem is to find the un-
known vector $\boldsymbol{\rho}$ from a family of quadratic equations

$$|\mathcal{A}_q\boldsymbol{\rho}|^2 = |\mathbf{b}_q|^2, \quad q = 1, \dots, Q,$$

470 understood component wise. This problem is nonlinear and nonconvex and, hence,
471 difficult to solve. In fact, it is in general NP hard [33]. However, if an appropriate set
472 of illuminations is used, we can take advantage of the polarization identity

$$473 \quad 2 \operatorname{Re} \langle u, v \rangle = |u + v|^2 - |u|^2 - |v|^2 \\ 474 \quad (3.16) \quad 2 \operatorname{Im} \langle u, v \rangle = |u - iv|^2 - |u|^2 - |v|^2$$

475 to solve a simple linear system of the form

476 (3.17)
$$\mathcal{A}_q \boldsymbol{\rho} = \mathbf{m}_q^{(r)}.$$

477 The polarization identity allows us to find the inner product between two complex
 478 numbers and, therefore, its phase differences. In (3.17), $\mathbf{m}_q^{(r)}$ is the vector whose
 479 i th component is the correlation $\overline{b_q^{(r)}} b_{\hat{\mathbf{e}}_i}^{(r)}$ between two signals measured at $\vec{\mathbf{x}}_r$, one
 480 corresponding to a general illumination $\hat{\mathbf{f}}_q(\omega)$ and the other to an illumination $\hat{\mathbf{e}}_i =$
 481 $[0, 0, \dots, 0, 1, 0, \dots, 0]^T$ whose entries are all zero except the i th entry which is 1.
 482 Using the polarization identity (3.16) we can obtain $\overline{b_q^{(r)}} b_{\hat{\mathbf{e}}_i}^{(r)}$ from linear combinations
 483 of the magnitudes (squared) $|b_q^{(r)}|^2$, $|b_{\hat{\mathbf{e}}_i}^{(r)}|^2$, $|b_q^{(r)} + b_{\hat{\mathbf{e}}_i}^{(r)}|^2$, and $|b_q^{(r)} + i b_{\hat{\mathbf{e}}_i}^{(r)}|^2$. A physical
 484 interpretation of (3.17) is as follows. Send an illumination $\hat{\mathbf{f}}_q(\omega)$, collect the response
 485 at $\vec{\mathbf{x}}_r$, time reverse the received signal at $\vec{\mathbf{x}}_r$, and send it back to probe the medium
 486 again. Then, $\mathbf{m}_q^{(r)}$ represents the signals recorded at all receivers $\vec{\mathbf{x}}_i$, $i = 1, \dots, N$.

487 To wrap up, if the phases are not measured at the array but we control the
 488 illuminations, the images can be formed by solving (3.17). We can use ℓ_1 -norm
 489 minimization if only one vector $\mathbf{m}_q^{(r)}$ is obtained in the data acquisition process,
 490 or we can use MUSIC if enough vectors of this form are available [31, 28]. Note that
 491 in this approach, where only one frequency ω is used, the receiver $\vec{\mathbf{x}}_r$ is fixed.

492 **3.4. Multiple frequency signals and multiple receivers.** Finally, we con-
 493 sider the most general case in which multiple frequency signals are used to probe
 494 the medium from several source positions, and the echoes are measured at several
 495 receiver positions. This case considers all the possible diversity of information that
 496 can be obtained from the illuminations. We discuss first the situation in which the
 497 receivers measure amplitudes and phases and, then, the situation in which they can
 498 only measure amplitudes squared.

499 **3.4.1. Imaging with phases.** Assume that the data (including phases)

500 (3.18)
$$d(\vec{\mathbf{x}}_r, \vec{\mathbf{x}}_s, \omega_l) = P(\vec{\mathbf{x}}_r, \vec{\mathbf{x}}_s; \omega_l),$$

501 for all receiver locations $\vec{\mathbf{x}}_r$, source locations $\vec{\mathbf{x}}_s$, and frequencies ω_l are available for
 502 imaging. For an array with N colocated sources and receivers that emit S differ-
 503 ent frequencies the number of measurements is then equal to $N^2 S$. To make use of
 504 the coherence of these data over all the frequencies we could stack them in a col-
 505 umn vector \mathbf{b} , but then we would have to deal with a huge linear system $\mathcal{A} \boldsymbol{\rho} = \mathbf{b}$
 506 of size $N^2 S \times K$. To reduce the number of data used in an ℓ_1 approach, we con-
 507 sider that the illumination is of separable form, i.e., $\hat{\mathbf{f}}(\omega_l) = f(\omega_l) \mathbf{f}$ and the same
 508 vector \mathbf{f} is used for all the frequencies ω_l , $l = 1, \dots, S$. Thus, for an illumination
 509 $\hat{\mathbf{f}} = [\hat{\mathbf{f}}(\omega_1)^T, \hat{\mathbf{f}}(\omega_2)^T, \dots, \hat{\mathbf{f}}(\omega_S)^T]^T$ we stack the data (including phases) in a column
 510 vector

511 (3.19)
$$\mathbf{b}_{\hat{\mathbf{f}}} = [\mathbf{b}_{\hat{\mathbf{f}}(\omega_1)}^T, \mathbf{b}_{\hat{\mathbf{f}}(\omega_2)}^T, \dots, \mathbf{b}_{\hat{\mathbf{f}}(\omega_S)}^T]^T,$$

512 and we solve the system of equations

513 (3.20)
$$\mathcal{A}_{\hat{\mathbf{f}}} \boldsymbol{\rho} = \mathbf{b}_{\hat{\mathbf{f}}},$$

514 with the $(N \cdot S) \times K$ matrix

$$515 \quad (3.21) \quad \mathcal{A}_{\hat{f}} = \begin{pmatrix} \hat{g}_{\hat{f}(\omega_1)}^{(1)} \hat{\mathbf{g}}(\vec{\mathbf{y}}_1; \omega_1) & \hat{g}_{\hat{f}(\omega_1)}^{(2)} \hat{\mathbf{g}}(\vec{\mathbf{y}}_2; \omega_1) & \dots & \hat{g}_{\hat{f}(\omega_1)}^{(K)} \hat{\mathbf{g}}(\vec{\mathbf{y}}_K; \omega_1) \\ \hat{g}_{\hat{f}(\omega_2)}^{(1)} \hat{\mathbf{g}}(\vec{\mathbf{y}}_1; \omega_2) & \hat{g}_{\hat{f}(\omega_2)}^{(2)} \hat{\mathbf{g}}(\vec{\mathbf{y}}_2; \omega_2) & \dots & \hat{g}_{\hat{f}(\omega_2)}^{(K)} \hat{\mathbf{g}}(\vec{\mathbf{y}}_K; \omega_2) \\ \vdots & \vdots & \ddots & \vdots \\ \hat{g}_{\hat{f}(\omega_S)}^{(1)} \hat{\mathbf{g}}(\vec{\mathbf{y}}_1; \omega_S) & \hat{g}_{\hat{f}(\omega_S)}^{(2)} \hat{\mathbf{g}}(\vec{\mathbf{y}}_2; \omega_S) & \dots & \hat{g}_{\hat{f}(\omega_S)}^{(K)} \hat{\mathbf{g}}(\vec{\mathbf{y}}_K; \omega_S) \end{pmatrix}.$$

516 Here, $\hat{g}_{\hat{f}(\omega_l)}^{(j)} = \hat{\mathbf{g}}(\vec{\mathbf{y}}_j; \omega_l)^T \hat{\mathbf{f}}(\omega_l)$ denotes the field with frequency ω_l at position $\vec{\mathbf{y}}_j$.
 517 The system (3.20) relates the unknown vector $\boldsymbol{\rho} \in \mathbb{C}^K$ to the data vector $\mathbf{b}_{\hat{f}} \in \mathbb{C}^{(N \cdot S)}$
 518 in a coherent way. The system of linear equations (3.20) can, of course, be solved by
 519 appropriate ℓ_2 and ℓ_1 methods.

520 However, because (3.20) cannot be written in the form of an MMV problem, MU-
 521 SIC cannot be used to identify the support of $\boldsymbol{\rho}$ as in the previous imaging problems.
 522 The issue here is that matrix (3.21) cannot be factorized in the form $\mathcal{A}_{\hat{f}} = \tilde{\mathbf{A}} \Lambda_{\hat{f}}$
 523 because the scalars $\hat{g}_{\hat{f}(\omega_l)}^{(j)}$ depend on frequency. However, in the paraxial regime,
 524 where the scatterers are far from the array, and the array and the IW are small so the
 525 wavefronts that illuminate the scatterers are planar, we can take into account these
 526 changes over frequencies explicitly to image coherently with MUSIC.

527 Indeed, assume for simplicity that only one source at $\vec{\mathbf{x}}_s = (\mathbf{x}_s, 0)$ with cross-
 528 range vector $\mathbf{x}_s = (x_{sx}, x_{sy})$ emits the signals, i.e., for all the frequencies ω_l we use
 529 the N-vector $\hat{\mathbf{f}}(\omega_l) \equiv \hat{\mathbf{f}}_{l,s} = [0, 0, \dots, 0, 1, 0, \dots, 0]^T$ with all the entries equal to zero
 530 except the s th entry which is one. In the paraxial regime, where $\lambda \ll a \ll L$ and
 531 the IW is small compared to L , the illumination at position $\vec{\mathbf{y}}_j = (\mathbf{y}_j, L + \eta_j)$ can
 532 be approximated by $\hat{g}_{\hat{f}_{l,s}}^{(j)} \approx e^{i\kappa_l(\eta_j + (\mathbf{x}_s - \mathbf{y}_j)^2/2L)} \approx e^{i\kappa_l \eta_j} e^{i\kappa_c(\mathbf{x}_s - \mathbf{y}_j)^2/2L}$ and, thus,
 533 $\mathcal{A}_{\hat{f}_{l,s}} \approx \tilde{\mathbf{A}} \Lambda_{\hat{f}_{c,s}}$ where

$$534 \quad (3.22) \quad \tilde{\mathbf{A}} = \begin{pmatrix} \hat{\mathbf{h}}(\vec{\mathbf{y}}_1; \omega_1) & \hat{\mathbf{h}}(\vec{\mathbf{y}}_2; \omega_1) & \dots & \hat{\mathbf{h}}(\vec{\mathbf{y}}_K; \omega_1) \\ \hat{\mathbf{h}}(\vec{\mathbf{y}}_1; \omega_2) & \hat{\mathbf{h}}(\vec{\mathbf{y}}_2; \omega_2) & \dots & \hat{\mathbf{h}}(\vec{\mathbf{y}}_K; \omega_2) \\ \vdots & \vdots & \ddots & \vdots \\ \hat{\mathbf{h}}(\vec{\mathbf{y}}_1; \omega_S) & \hat{\mathbf{h}}(\vec{\mathbf{y}}_2; \omega_S) & \dots & \hat{\mathbf{h}}(\vec{\mathbf{y}}_K; \omega_S) \end{pmatrix}$$

535 with $\widehat{\mathbf{h}}(\vec{\mathbf{y}}_j; \omega_l) = e^{i\kappa_l \eta_j} \widehat{\mathbf{g}}(\vec{\mathbf{y}}_j; \omega_l)$, and

$$536 \quad (3.23) \quad \Lambda_{\widehat{f}_{c,s}} = \begin{pmatrix} e^{i\kappa_c(\mathbf{x}_s - \mathbf{y}_1)^2/2L} & 0 & & & \\ 0 & e^{i\kappa_c(\mathbf{x}_s - \mathbf{y}_2)^2/2L} & & & \\ & & \ddots & & \\ & & & 0 & e^{i\kappa_c(\mathbf{x}_s - \mathbf{y}_K)^2/2L} \end{pmatrix}.$$

537 In this approximation, the nonzero entries of the diagonal matrix (3.23) are given by
 538 the illumination relative to the central frequency κ_c . Then, the multiple-frequency
 539 MUSIC formulation is of the MMV form

$$540 \quad (3.24) \quad \widetilde{\mathcal{A}} \Lambda_{\widehat{f}_{c,s}} \boldsymbol{\rho} = B,$$

541 with $\widetilde{\mathcal{A}}$ as in (3.22), $\Lambda_{\widehat{f}_{c,s}}$ as in (3.23), and the $(N \cdot S) \times N$ matrix

$$542 \quad (3.25) \quad B = P^c = [P(\omega_1)^T, P(\omega_2)^T, \dots, P(\omega_S)^T]^T$$

543 corresponding to stacking the array response data matrices (3.4) for multiple frequen-
 544 cies in a column. With this data structure, multiple-frequency imaging can be carried
 545 out coherently using MUSIC with the column vectors of (3.22) as the imaging vectors.

546 We could have used instead the alternative data structure

$$547 \quad (3.26) \quad B = P^d = \begin{pmatrix} P(\omega_1) & \dots & 0 & 0 \\ 0 & P(\omega_2) & \dots & 0 \\ \dots & \dots & \dots & \dots \\ 0 & 0 & 0 & P(\omega_S) \end{pmatrix}$$

548 to image with MUSIC. However, that would be as if imaging with each frequency
 549 separately and summing up the resulting images incoherently, so there would be no
 550 significant improvement over single frequency imaging.

551 To summarize, multiple frequency imaging with phases can be done in all regimes
 552 by solving (3.20) with suitable ℓ_2 -norm or ℓ_1 -norm methods. The matrix-matrix
 553 formulation (3.24) can be used to form the images with MUSIC or using (2,1)-matrix
 554 minimization as in [12]. Recall that (3.24) is an approximate formulation, which is
 555 valid for the paraxial regime.

556 **3.4.2. Imaging without phases.** Assume now that only the intensities can be
 557 recorded at the array. In subsection 3.3 we showed that with multiple sources and
 558 multiple receivers, but a single frequency, we could recover cross correlated data from
 559 intensity-only measurements if we control the illuminations and, then, we could image
 560 holographically. In general, if several frequencies are used for imaging, we can fix one
 561 of the three possible variables $(\vec{\mathbf{x}}_r, \vec{\mathbf{x}}_s, \omega)$ and proceed similarly. For example, we can
 562 fix the receiver position $\vec{\mathbf{x}}_r$, and recover the multifrequency interferometric data

$$563 \quad (3.27) \quad d((\vec{\mathbf{x}}_r, \vec{\mathbf{x}}_r), (\vec{\mathbf{x}}_s, \vec{\mathbf{x}}_{s'}), (\omega, \omega')) = \overline{P(\vec{\mathbf{x}}_r, \vec{\mathbf{x}}_s; \omega)} P(\vec{\mathbf{x}}_r, \vec{\mathbf{x}}_{s'}; \omega')$$

564 for all pairs of frequencies (ω, ω') and source locations $(\vec{\mathbf{x}}_s, \vec{\mathbf{x}}_{s'})$.

565 To understand the type of data that we can use in this situation, let us consider
 566 one row of the $N \times (N \cdot S)$ full response matrix for multiple frequencies

$$567 \quad (3.28) \quad P^r = [P(\omega_1), P(\omega_2), \dots, P(\omega_S)],$$

568 and denote the r -th row of this matrix by

$$569 \quad (3.29) \quad \mathbf{p}_r = [p_{r1}, p_{r2}, \dots, p_{rN \cdot S}].$$

570 Here, p_{rj} with $j \equiv j(s, l) = s + (l - 1) \cdot N$, denotes the received signal at $\vec{\mathbf{x}}_r$ when the
571 source at $\vec{\mathbf{x}}_s$ sends a signal of frequency ω_l . With this notation, and denoting by the
572 superscript \cdot^* the conjugate transpose of a vector,

$$573 \quad (3.30) \quad M^r = \mathbf{p}_r^* \mathbf{p}_r$$

574 is the rank-one matrix whose j th column corresponds to the vector $\mathbf{m}_{\hat{\mathbf{e}}_j}^r$ in the right
575 hand side of the linear system (3.17), introduced in subsection 3.3 for single frequency
576 imaging, but generalized here so as to account for multiple frequencies, i.e., for $l =$
577 $1, \dots, S$. That is, the j th column of (3.30) contains the correlations of the response
578 received at $\vec{\mathbf{x}}_r$ when a signal of unit amplitude and frequency ω_l is sent from $\vec{\mathbf{x}}_s$ to
579 probe the medium ($j = s + (l - 1)N$), with all the other responses received also at
580 $\vec{\mathbf{x}}_r$ when unit amplitude signals are sent from all the sources with all the different
581 frequencies. In short,

$$582 \quad (3.31) \quad [M^r]_{ij} = \bar{p}_{ri} p_{rj} = (\mathbf{p}_r \hat{\mathbf{e}}_i)^* \mathbf{p}_r \hat{\mathbf{e}}_j.$$

583 Since M^r is rank one, all the columns are linearly dependent, so we can only use
584 one of its columns to solve the imaging problem

$$585 \quad (3.32) \quad \mathcal{A}_{\hat{\mathbf{e}}_j} \boldsymbol{\rho} = \mathbf{m}_{\hat{\mathbf{e}}_j}^r$$

586 for one $\hat{\mathbf{e}}_j$, and form the images with an ℓ_2 -norm or ℓ_1 -norm method. The matrix $\mathcal{A}_{\hat{\mathbf{e}}_j}$
587 is given by (3.21) and, hence, the model (3.32) is exact.

588 Alternatively, once the matrix M^r has been obtained from intensity-only mea-
589 surements, imaging can be done using the Kirchhoff migration functional

$$590 \quad (3.33) \quad \mathcal{I}^{\text{KM}} = \text{diag}(\mathcal{A}_{\hat{\mathbf{e}}_j}^* M^r \mathcal{A}_{\hat{\mathbf{e}}_j}).$$

591 The ℓ_2 images (3.33) are very robust with respect to additive measurement noise, but
592 they are statistically unstable when imaging is done in a randomly inhomogeneous
593 medium or when there are modeling errors due to off-grid scatterers. Both situations
594 lead to perturbations in the (unknown) phases that may make the \mathcal{I}^{KM} images depen-
595 dent on the particular realization of the medium and/or the positions of the scatterers.
596 In [29], we showed that statistical stability can be enhanced by masks that limit the
597 frequency and source offsets of the measurements used in (3.33). Hence, if the pertur-
598 bations of the phases are important, we can use the Single Receiver INTerferometric
599 (SRINT) imaging functional given by

$$600 \quad (3.34) \quad \mathcal{I}^{\text{SRINT}} = \text{diag}(\mathcal{A}_{\hat{\mathbf{e}}_j}^* \mathcal{Z} \odot M^r \mathcal{A}_{\hat{\mathbf{e}}_j}).$$

601 In (3.34), the mask \mathcal{Z} is a matrix composed by zeros and ones restricting the data
602 to coherent nearby source locations and frequencies, and \odot denotes component-wise
603 multiplication. The same idea can be used for stabilizing the ℓ_1 -norm minimization
604 method if the perturbation of the phases are important. We can just replace the j th
605 column of the matrix M^r by the j th column of the masked data $\mathcal{Z} \odot M^r$, and remove
606 the corresponding rows from the model matrix $\mathcal{A}_{\hat{\mathbf{e}}_j}$.

607 On the other hand, as noted in [31, 28], the support of the reflectivity $\boldsymbol{\rho}$ can be
608 recovered exactly by using the MUSIC algorithm on the single frequency interfero-
609 metric matrix $M(\omega) = P^*(\omega)P(\omega)$. Once the support of $\boldsymbol{\rho}$ is found, we can estimate

610 the reflectivities by solving a trace minimization problem restricted to the support of
 611 $\boldsymbol{\rho}$ (see [10, 31] for details).

612 For multiple frequencies, multiple sources and multiple receivers one can use the
 613 data structure

$$614 \quad (3.35) \quad M^c = \begin{pmatrix} P(\omega_1)^*P(\omega_1) \\ P(\omega_2)^*P(\omega_1) \\ \vdots \\ P(\omega_S)^*P(\omega_1) \end{pmatrix}$$

615 for pairs of frequencies (ω_l, ω_1) , $l = 1, \dots, S$, to image coherently using MUSIC.
 616 Indeed, the matrices M^c as in (3.35) and P^c defined in (3.25) have the same column
 617 space and, therefore, MUSIC can form the images using the SVD of M^c and the
 618 column vectors of (3.22) as imaging vectors. We denote these data structures with
 619 the superscript c to point out that we have stacked the one frequency matrices $P(\omega_l)$
 620 and the two frequencies matrices $P(\omega_l)^*P(\omega_1)$ in a column.

621 As noted in the previous section we could have used instead the alternative data
 622 structure

$$623 \quad (3.36) \quad M^d = \begin{pmatrix} P(\omega_1)^*P(\omega_1) & \dots & 0 & 0 \\ 0 & P(\omega_2)^*P(\omega_2) & \dots & 0 \\ \dots & \dots & \dots & \dots \\ 0 & 0 & 0 & P(\omega_S)^*P(\omega_S) \end{pmatrix}$$

624 to image using MUSIC. However, as we have already explained, if we used the SVD
 625 of M^d to obtain the signal and noise subspaces, then the frequencies are not used
 626 coherently and there is no improvement over single frequency imaging.

627 In summary, multiple frequency imaging with intensity-only can be done in all
 628 regimes by solving (3.32) with appropriate ℓ_2 -norm or ℓ_1 -norm methods or, in the
 629 paraxial regime, by forming the images using MUSIC on the data structure (3.35)
 630 with imaging vectors given by the column vectors of the matrix (3.22). MUSIC on the
 631 data structure (3.36) should not be used since multiple frequencies are not processed
 632 coherently. The performance of these methods will be assessed in Section 4, where we
 633 show numerical experiments in homogeneous and weakly inhomogeneous media.

634 **4. Numerical Simulations.** We present here numerical simulations that illus-
 635 trate the performance of the different imaging methods discussed in the previous sec-
 636 tions. Specifically, we consider multifrequency interferometric imaging without phases
 637 discussed in subsection 3.4.2, and we present the images obtained with ℓ_1 -norm min-
 638 imization, SRINT, and MUSIC using the data structures M^c and M^d . Our objective
 639 is to study the robustness of these imaging methods in the presence of noise, that
 640 is perturbations in the unknown phases of the collected signals. Two types of phase
 641 perturbations are considered, systematic due to off-grid placement of the scatterers
 642 and random resulting from wave propagation in an inhomogeneous ambient medium.

643 **4.1. Imaging setup.** We consider a typical imaging regime in optics, with a
 644 central frequency $f_0 = 600$ THz corresponding to a central wavelength $\lambda_0 = 500nm$.
 645 We use $S = 12$ equally spaced frequencies covering a total bandwidth of 30THz. In
 646 this regime, the decoherence frequency of the data Ω_d is equal to the total bandwidth.
 647 All considered wavelengths are in the visible spectrum of green light.

648 The size of the array is $a = 500\lambda_0$, and the distance between the array and the IW
 649 is $L = 10000\lambda_0$. The IW, whose size is $120\lambda_0 \times 60\lambda_0$, is discretized using a uniform

650 lattice with mesh size $4\lambda_0 \times 2\lambda_0$. The medium between the array and the IW is
 651 inhomogeneous, with weak fluctuations and long correlation lengths with respect to
 652 the central wavelength. The propagation distance L is large so cumulative scattering
 653 effects are important, but not too large so the phases of the signals received at the
 654 array still maintain certain degree of coherence. In all the figures, the true locations
 655 of the scatterers are indicated with white crosses, and the length scales are measured
 656 in units of λ_0 .

657 Again, we assume that the phases of the signals received at the array cannot be
 658 measured. Hence, only their intensities are available for imaging. These measure-
 659 ments are collected at only one receiver, so we can use the methods explained in
 660 subsection 3.4.2 to image interferometrically. We consider imaging in homogeneous
 and inhomogeneous media.

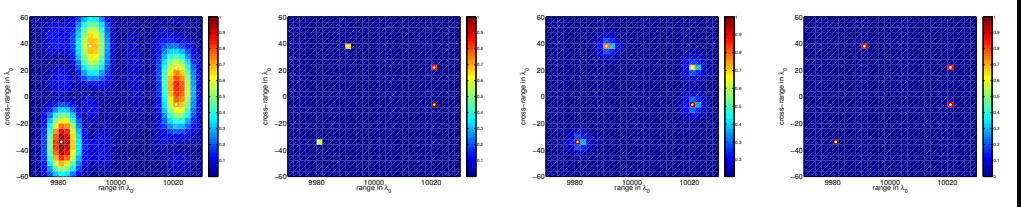


FIG. 2. *Imaging in a homogeneous medium. There is no noise added to the data and the scatterers are on the grid. From left to right: SRINT image, MUSIC with M^d , MUSIC with M^c coupling over frequencies, and ℓ_1 -norm minimization applied on one column of the masked matrix $\mathcal{Z} \odot M^r$.*

661

662 **4.2. Imaging in homogeneous media.** Let us first consider imaging in ho-
 663 mogeneous media. For the imaging system described above, we expect cross-range
 664 and range resolutions of $\lambda_0 L/a = 20\lambda_0$ and $C_0/B = \lambda_0 f_0/B = 20\lambda_0$, respectively. In
 665 order to keep the resolution fixed with respect to imaging in inhomogeneous media
 666 that we consider afterwards, we also apply masks to the data used to image in the
 667 homogeneous medium. This reduces the cross-range resolution to $\lambda_0 L/X_d = 32\lambda_0$
 668 corresponding to $X_d = 5a/8$. The range resolution does not change because the
 669 decoherence frequency Ω_d is equal to the total bandwidth.

670 In Figure 2, the scatterers lie on the grid and there is no noise in the data. We
 671 observe that SRINT (left image) provides a quite limited resolution and it cannot
 672 resolve two of the four scatterers. On the other hand, imaging with MUSIC (two
 673 middle images) or imaging using ℓ_1 -norm minimization (right image) give much better
 674 results. MUSIC using the block-diagonal matrix M^d (second image from the left) gives
 675 exact recovery, while MUSIC using the M^c matrix (third image from the left), that
 676 couples all the frequencies, is less accurate. This is so because, as we explained in
 677 Section 3.4, MUSIC with M^c is not exact as it provides approximate locations of the
 678 scatterers only in the paraxial regime. Finally, the ℓ_1 -norm approach recovers exactly
 679 the four scatterers as can be seen in the right image of this figure.

680 Figure 3 shows the same experiment as Figure 2 but with the scatterers displaced
 681 by half the grid size with respect to the grid points in range and cross-range directions.
 682 This produces perturbations in the unknown phases of the collected signals due to
 683 modeling errors. Because the point spread function is, in this case, much wider (of the
 684 order of $20\lambda_0$) than the off-grid displacements, the image formed with SRINT (left
 685 plot) is very robust with respect to these perturbations in the phases. However, the
 686 image obtained with MUSIC using the data structure M^d (second plot from the left)

687 deteriorates dramatically because the multiple-frequency information contained in the
 688 data is not processed in a coherent way. On the other hand, both MUSIC with the
 689 M^c data structure (third plot from the left) and ℓ_1 -norm minimization (right plot)
 690 are very robust with respect to the off-grid displacements.

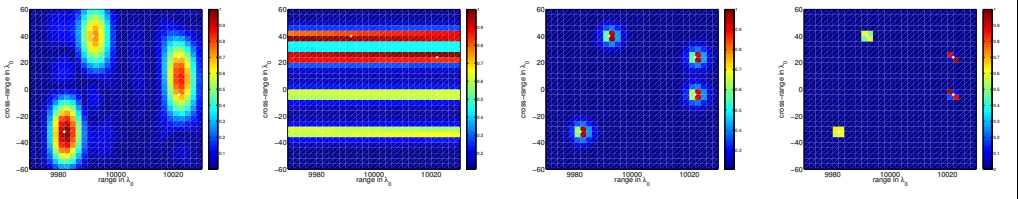


FIG. 3. Same as Figure 2 but with the scatterers off the grid. The scatterers are displaced by half the grid size in both directions from a grid point.

691 We study next the performance of the proposed methods for imaging in inhomogeneous media with weak fluctuations and long correlation lengths with respect to λ_0 .
 692 The challenge is to obtain similar results in this case.
 693

694 **4.3. Imaging in random inhomogeneous media.** Consider the setup displayed in Figure 4 with four scatterers in the right (black circles) at a distance $L = 10000\lambda_0$ from the array (black stars). The data used in the numerical experiments are generated using the random phase model which is frequently used to account for weak phase distortions [3, 13, 5, 29]. In this model, the standard deviation of the perturbations of the phases is given by $\sigma\sqrt{lL}/\lambda_0$, where σ and l denote the strength and the correlation length of the fluctuations of the medium, respectively. If we introduce the characteristic strength $\sigma_0 = \lambda_0/\sqrt{lL}$, for which the standard deviation of the random phases is $O(1)$, we can quantify the perturbations of the unknown phases by the dimensionless parameter $\varepsilon = \sigma/\sigma_0$.

704 In order to study the effect of phase distortions due to a random medium on imaging, we consider that the scatterers lie on the grid. Imaging in random media with ℓ_1 -norm minimization has also been considered in [13, 5].
 705
 706

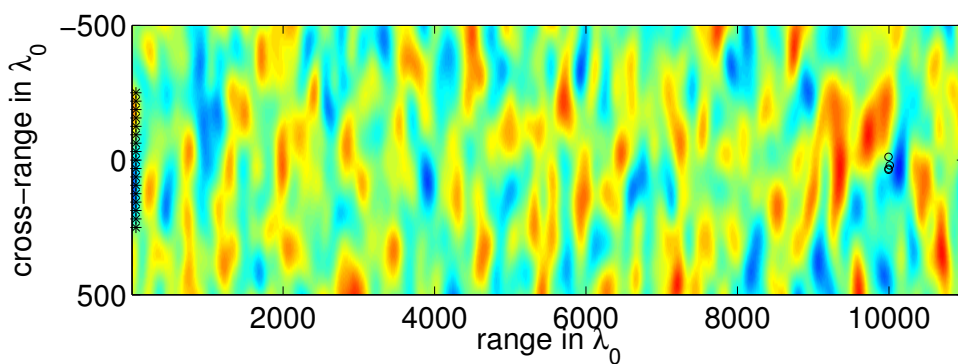


FIG. 4. One realization of the random medium used in the simulations. The correlation length of the fluctuations is $l = 100\lambda_0$.

707 Figure 5 displays the images obtained in a very weak fluctuating random medium with $\varepsilon = 0.05$. Comparing these images with the ones obtained in a homogeneous
 708

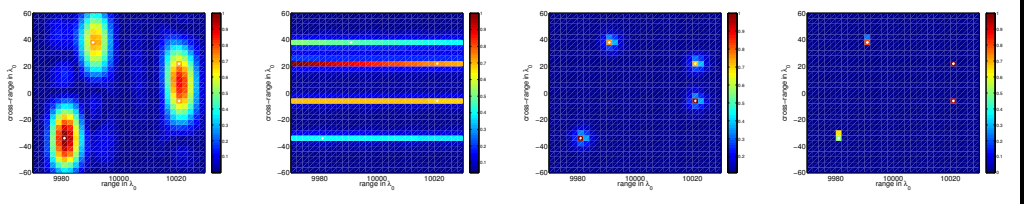


FIG. 5. Same as Figure 2 but the medium is inhomogeneous. The strength of the fluctuations is $\sigma = 0.5 \cdot 10^{-4}$ which corresponds to $\varepsilon = 0.05$. The scatterers are on-grid.

709 medium with scatterers on and off the grid (Figures 2 and 3, respectively) we observe
 710 that (i) SRINT (left plot), MUSIC using M^c (third plot from the left) and ℓ_1 -norm
 711 minimization (right plot) are stable, and (ii) MUSIC using M^d (second plot from the
 712 left) is not. Note that off-grid scatterers and a random medium both induce similar
 713 noise in the data, as both occur in the phases. In the off-grid case, the noise is
 714 systematic and similar for all array elements, while the noise induced by the random
 715 phase model depends on the path that connects the scatterer to each array element.
 716 Hence, depending on the correlation length of the random medium the noise produced
 717 in the phases is more or less correlated over the array elements.

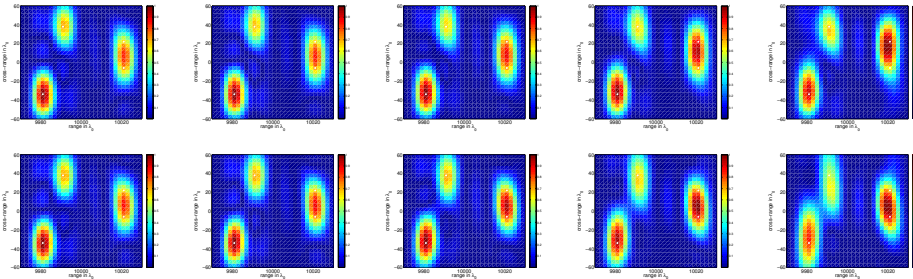


FIG. 6. Imaging with SRINT in inhomogeneous media illustrating its stability with respect to the random fluctuations of the media. The strength of the fluctuations increases from left to right so $\varepsilon = 0.1, 0.2, 0.4, 0.6$ and 0.8 . The top and bottom rows are two realizations of the random medium.

718 Since MUSIC using M^d is not robust with respect to perturbations in the phases
 719 (see Figures 3 and 5) because the data are not processed coherently over frequencies,
 720 we do not present more results using this method.

721 To further examine the robustness of the other imaging methods with respect to
 722 random medium fluctuations, we consider in the next figures five noise levels corre-
 723 sponding to $\varepsilon = 0.1, 0.2, 0.4, 0.6$ and 0.8 . Each figure presents results for two realiza-
 724 tions of the random medium. In Figure 6 we see that, as expected, SRINT is highly
 725 robust, although its resolution is not very good. Even for $\varepsilon = 0.8$ (right column) the
 726 images do not change much respect to the ones obtained in a homogeneous medium.
 727 Figure 7 shows the images obtained with ℓ_1 -norm minimization. The resolution is
 728 much better than that provided by SRINT, but it is much more sensitive to noise.
 729 Only for fluctuation strengths below or equal $\varepsilon = 0.2$ the images are good. Above this
 730 strength the images are useless. However, the use of masks on the data effectively
 731 removes the distortion imposed by the medium up to $\varepsilon = 0.4$, as it can be seen in
 732 Figure 8. This is so because by using masks we discard the incoherent data and, thus,
 733 we improve the robustness of the ℓ_1 -norm method (even though we reduce the number

734 of equations in the linear system by about 40%).

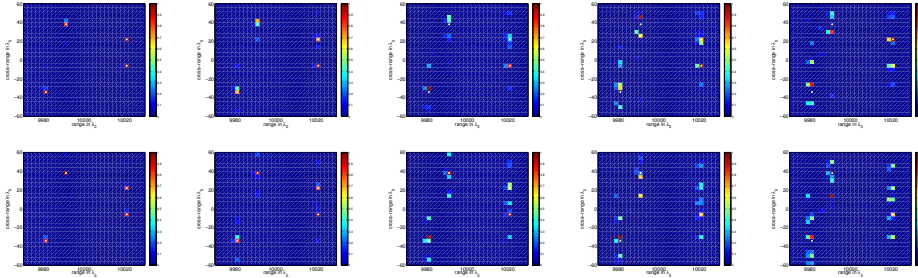


FIG. 7. Images obtained with ℓ_1 -norm minimization without masks in the same media and the same scatterer's configuration as in Figure 6. Imaging with ℓ_1 -norm minimization without masks is stable only for $\varepsilon \leq 0.2$.

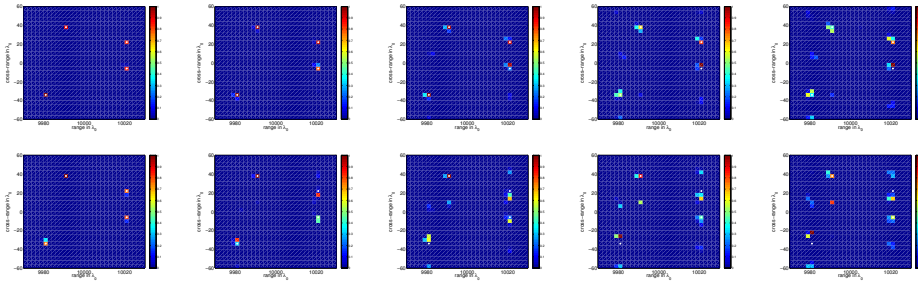


FIG. 8. Same as Figure 7 but using masked data. The results are now stable for $\varepsilon \leq 0.4$.

735 Finally, the images shown in Figure 9 formed using MUSIC with M^c are also very
 736 good. They have significantly better resolution than the SRINT images but not as
 737 good as the ones obtained with ℓ_1 -norm minimization. We stress that MUSIC with
 738 M^c is not exact even for perfect data and, therefore, ℓ_1 -norm minimization should
 739 be preferred if the fluctuations of the medium are weak. However, as the strength
 740 of the fluctuations increases, MUSIC with M^c becomes competitive. Observe that at
 741 lower SNR, when the ℓ_1 -norm images are not useful, MUSIC with M^c is robust and
 742 the resolution is better than the one provided by SRINT. Therefore, it should be the
 743 preferred method among the three for imaging in moderate SNR regimes.

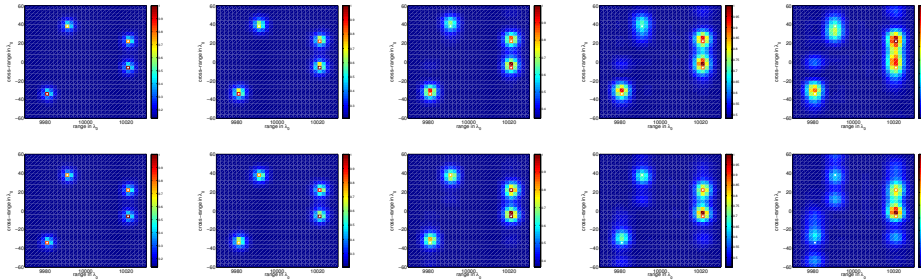


FIG. 9. Images obtained with MUSIC using M^c in the same media and the same scatterer's configuration as in Figures 6-8. MUSIC using M^c is stable for $\varepsilon \leq 0.6$.

5. Imaging results in the framework of Theorems 2.3 and 2.4. To illustrate the relevance of Theorems 2.3 and 2.4 for imaging, we consider in this section the equivalent source problem of active array imaging with multiple frequencies and multiple receivers described in subsection 3.4.1. In this setting we have to solve the linear system

$$\tilde{\mathcal{A}}\boldsymbol{\rho} = \mathbf{b}_{\tilde{\mathcal{f}}}$$

with $\tilde{\mathcal{A}}$ the model matrix (3.22). We compare the corresponding ℓ_2 and ℓ_1 solutions of this problem for different imaging configurations. Our results illustrate the well known super-resolution for ℓ_1 , meaning that $\boldsymbol{\rho}_{\ell_1}$ determines the support of the unknown $\boldsymbol{\rho}$ with higher accuracy than the conventional resolution limits, provided the assumptions of Theorem 2.3 for the noiseless case or Theorem 2.4 for the noisy case are satisfied. We also show how the bandwidth, the array size and the number of scatterers affect the vicinities defined in (2.8). The numerical results are not specialized to a particular physical regime. They illustrate only the role of the theorems in solving the associated linear systems.

Imaging methods. We compare the solution $\boldsymbol{\rho}_{\ell_1}$ obtained with the ℓ_1 -norm minimization algorithm GelMa described in section 2, and the ℓ_2 -norm solution

$$(5.1) \quad \boldsymbol{\rho}_{\ell_2} = \tilde{\mathcal{A}}^* \mathbf{b}_{\tilde{\mathcal{f}}}.$$

where $\tilde{\mathcal{A}}^*$ is the conjugate transpose of $\tilde{\mathcal{A}}$.

Imaging setup. The images are obtained in a homogeneous medium with an active array of $N = 37$ transducers. The ratio between the array aperture a and the distance L to IW, as well as the ratio between the bandwidth $2B$ and the central frequency f_0 , vary in the numerical experiments. The IW is discretized using a uniform grid of $K = 3721$ points of size $\lambda_0/2$ in range and cross-range directions. The classical resolution theory suggests that the range and cross-range resolutions are $c_0/(2B)$ and $\lambda_0 L/a$, respectively. There is no additive noise in the data, but we consider on-grid and off-grid scatterers which produces perturbations in the recorded phases.

Imaging results. In Figure 10 we show the results obtained for a large array and a large bandwidth corresponding to $a/L = 1$ and $(2B)/f_0 = 1$. From left to right we show the $\boldsymbol{\rho}_{\ell_2}$ solution, the $\boldsymbol{\rho}_{\ell_1}$ solution, and the vicinities S_j defined in (2.8) plotted with different colors. In the top and bottom rows there are $M = 4$ and $M = 8$ scatterers, respectively. All the scatterers are on the grid and their exact locations are indicated with white crosses. The four scatterers in the top row are far apart and,

771 therefore, their vicinities do not overlap as it can be seen in the top right image of
772 this figure. In this case, all the conditions of Theorem 2.3 are satisfied and we find
773 the exact locations of scatterers with the ℓ_1 -norm minimization algorithm. The eight
774 scatterers in the bottom row are closer and their vicinities are larger (according to
775 (2.8) the size of the vicinities increases with M). We observe in the bottom right image
776 of this figure that the vicinities overlap, so condition (2.10) is not satisfied in this case.
777 We still, however, find the exact locations of scatterers with the ℓ_1 -norm minimization
778 algorithm which means that the conditions of Theorem 2.3 have pessimistic bounds.
779 Because the array and the bandwidth are large, the ℓ_2 -norm solutions also give very
good estimates of the scatterer's locations (see the left column images).

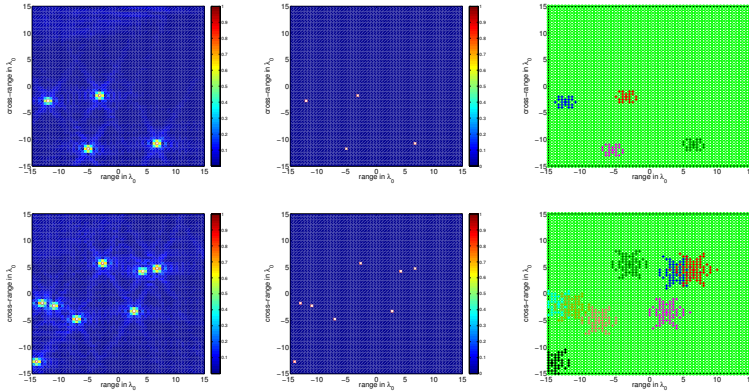


FIG. 10. *Imaging in a homogeneous medium and scatterers on grid. From left to right: ρ_{ℓ_2} , ρ_{ℓ_1} , and the vicinities S_j , $j = 1, \dots, M$, plotted with different colours. Top row $M = 4$, bottom row $M = 8$. Large array aperture and large bandwidth so $a/L = 1$ and $(2B)/f_0 = 1$.*

780

781 In Figure 11 we show the results for the same configurations of scatterers as in
782 Figure 10, but using a smaller array aperture and a smaller bandwidth so $a/L = 1/2$
783 and $(2B)/f_0 = 1/2$. Thus, the classical resolution limits become $c_0/(2B) = 2\lambda_0$ in
784 range and $\lambda_0 L/a = 2\lambda_0$ in cross-range. Hence, the resolution of the ℓ_2 -norm solutions
785 deteriorate, as can be observed in the left column images of this figure. In fact, we only
786 recover seven scatterers instead of eight for $M = 8$ (there are two scatterers that are
787 quite close). The ℓ_1 -norm minimization approach, however, still gives exact recovery
788 for both $M = 4$ and $M = 8$ scatterers. This is referred to as super-resolution, which
789 means that we can determine the location of the scatterers with a better accuracy
790 than the classical resolution limits.

791 To illustrate the effect of the array and bandwidth sizes on the size of the vicinities
792 we plot them in Figure 12 for the case $M = 4$. From left to right we plot the vicinities
793 for $a/L = 1/2$ and $(2B)/f_0 = 1/2$, $a/L = 1/2$ and $(2B)/f_0 = 1/4$, and $a/L = 1/4$
794 and $(2B)/f_0 = 1/2$. As expected, cross-range and range resolutions deteriorate and
795 consequently vicinity sizes increase as the ratios a/L and $(2B)/f_0$ decrease.

796 In Figure 13 we use a relatively small array and bandwidth so $a/L = 1/4$ and
797 $(2B)/f_0 = 1/4$. In this case, the conditions of Theorem 2.3 are not satisfied for neither
798 $M = 4$ nor $M = 8$, but the images obtained with ℓ_1 -norm minimization are still very
799 good. They are exact for $M = 4$ and very close to the true image for $M = 8$.

800 By further decreasing the array aperture and the bandwidth so that $a/L = 0.1$
801 and $(2B)/f_0 = 0.1$, we consider in Figure 14 a very challenging situation even for well
802 separated scatterers. The ℓ_2 -norm solutions shown in the left column of this figure

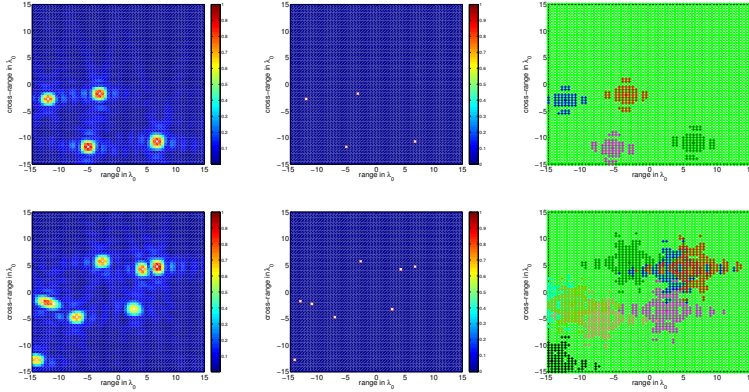


FIG. 11. Same as Figure 10 but using a smaller array aperture and a smaller bandwidth so $a/L = 1/2$ and $(2B)/f_0 = 1/2$.

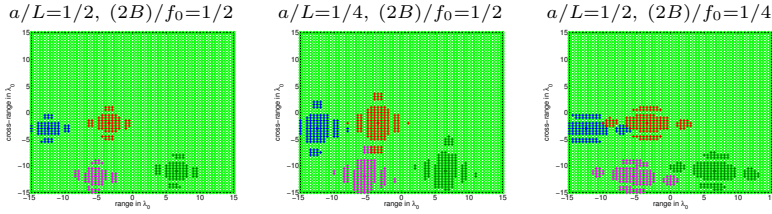


FIG. 12. Vicinities S_j , $j = 1, \dots, 4$, for different array and bandwidth sizes. From left to right: $a/L = 1/2$ and $(2B)/f_0 = 1/2$, $a/L = 1/2$ and $(2B)/f_0 = 1/4$ and $a/L = 1/4$ and $(2B)/f_0 = 1/2$.

803 are not able to locate the positions of the scatterers because of the low resolution of
 804 the imaging system. However, when the number of the scatterers is very small (see
 805 the top row corresponding to $M = 4$) the ℓ_1 -norm approach provides a precise image
 806 even though the discretization of the IW is 20 times finer than the classical resolution
 807 limits of the imaging system. On the other hand, when we increase the number of
 808 scatterers to $M = 8$ (bottom row) the interaction between the vicinities is very strong
 809 and the ℓ_1 -norm image is not good neither.

810 We now consider the same situation as in Figure 10, so the array aperture and
 811 the bandwidth are large, but with scatterers off the grid. This means that there are
 812 modeling errors and, therefore, there is not a vector ρ for which $\tilde{\mathcal{A}}\rho = \mathbf{b}_{\hat{f}}$. In the
 813 case considered next, the scatterers are displaced by $\lambda_0/4$ from a grid point in range
 814 and cross-range directions. The left column of Figure 15 shows, as expected, that the
 815 ℓ_2 -norm solutions (5.1) are not affected by off-grid displacements. This is so because
 816 the resolution is larger than the displacements of the scatterers with respect to the
 817 grid points. The right column shows, however, that the ℓ_1 -norm solutions are sensitive
 818 to these displacements. They are no longer exact, although they remain very close to
 819 the true solutions. By carefully examining the results of this figure we observe that
 820 the ℓ_1 -norm solutions behave as it is predicted by Theorem 2.4. The coherent part of
 821 the solution is supported in the vicinities of the exact solution while the incoherent
 822 part remains very small.

823 Figure 16 shows similar results but for a smaller array and a smaller bandwidth.
 824 We use $a/L = 1/4$ and $(2B)/f_0 = 1/4$, so the classical resolution limits increase as

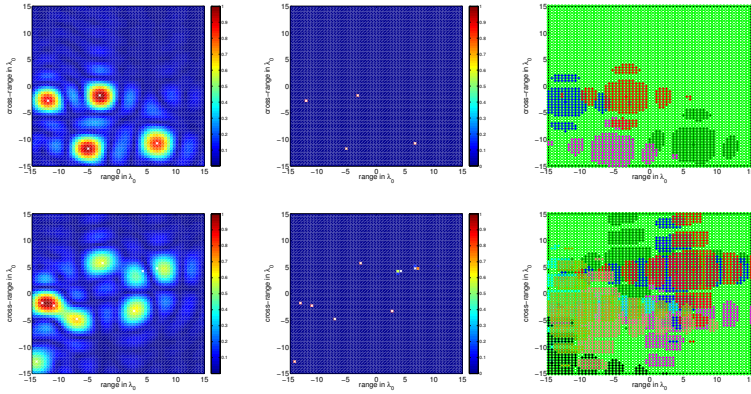


FIG. 13. Same as Figures 10 and 11 but using a smaller array aperture and a smaller bandwidth so $a/L = 1/4$ and $(2B)/f_0 = 1/4$.

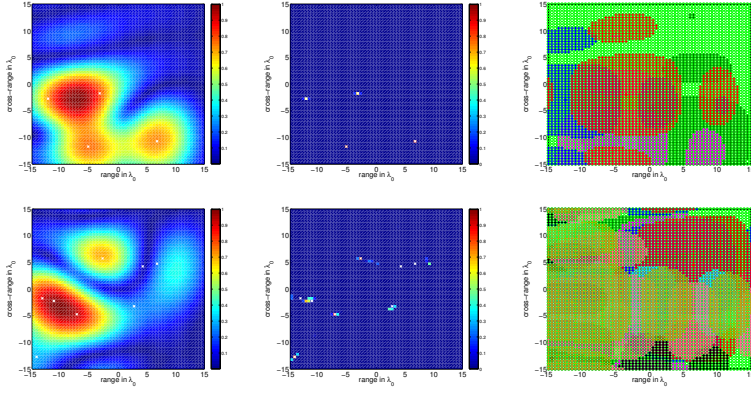


FIG. 14. Imaging in a homogeneous medium with $a/L = 0.1$ and $(2B)/f_0 = 0.1$. Top and bottom rows: $M = 4$ and $M = 8$ scatterers, respectively. From left to right: ρ_{ℓ_2} as in (5.1), ρ_{ℓ_1} obtained with GelMa, and the vicinities S_j , $j = 1, \dots, M$ plotted with different colors.

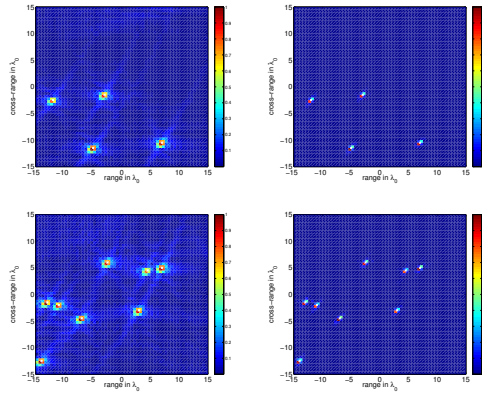


FIG. 15. Imaging in a homogeneous medium with scatterers off the grid. As in Figure 10, we use a large array aperture and a large bandwidth so $a/L = 1$ and $(2B)/f_0 = 1$. Top and bottom rows show the images for $M = 4$ and $M = 8$ scatterers, respectively. Left and right columns show the ℓ_2 -norm and ℓ_1 -norm solutions, respectively.

825 can be observed in the ℓ_2 -norm solutions shown in the left column. As in the previous
 826 figure, the ℓ_1 -norm solutions shown in the right column have a coherent part whose
 827 support is contained in the vicinities of the true solutions and an incoherent part that
 828 is very small. We also refer to [18, 4] for nice discussions about what to expect from
 ℓ_1 -norm minimization when the scatterers do not lie on the grid.

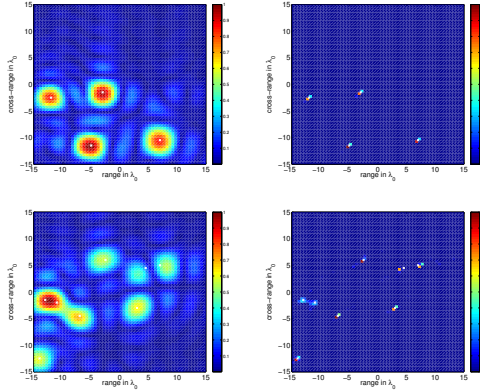


FIG. 16. Same as Figure 15 but with $a/L = 1/4$ and $(2B)/f_0 = 1/4$.

829

830 **6. Conclusions.** In this paper we addressed the question of what are appro-
 831 priate data structures so as to obtain robust images with two widely used methods:
 832 ℓ_1 -norm minimization and MUSIC. Both methods are well adapted to finding sparse
 833 solutions of linear underdetermined systems of equations of the form $\mathcal{A}_l \boldsymbol{\rho} = \mathbf{b}_l$ where
 834 \mathbf{l} is a parameter vector that can be varied, such as the illumination profile in space
 835 and/or frequency. ℓ_1 -norm minimization is well suited for solving problems with a
 836 single measurement vector corresponding to one parameter vector \mathbf{l} . On the other
 837 hand, MUSIC requires multiple measurement vectors that are obtained for several
 838 parameter vectors \mathbf{l}_i , $i = 1, \dots, S$. Given the data \mathbf{b}_l , our first main result concerns
 839 the uniqueness and robustness to noise of the minimal ℓ_1 -norm solution of $\mathcal{A}_l \boldsymbol{\rho} = \mathbf{b}_l$.
 840 This is the subject of Theorems 2.3 and 2.4. The second important result is the key
 841 observation that MUSIC provides the exact support of the unknown $\boldsymbol{\rho}$ when the ma-
 842 trix \mathcal{A}_l admits a factorization of the form $\mathcal{A}_l = \tilde{\mathcal{A}} \Lambda_l$ with Λ_l diagonal. Furthermore,
 843 we show in Theorem 2.6 that MUSIC is robust with respect to noise. Our third main
 844 contribution is the formulation of several common imaging configurations, including
 845 multifrequency imaging and imaging without phases, under a common linear algebra
 846 framework. For imaging without phases (the phase retrieval problem) the robustness
 847 of ℓ_1 -norm minimization and MUSIC is studied with numerical simulations in weakly
 848 inhomogeneous media. Our results suggest that ℓ_1 -norm minimization may be used
 849 for low noise levels while MUSIC should be the method of choice for higher noise
 850 levels.

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946 Appendix A. Proofs of theorems 2.1 to 2.4.

947 THEOREM 2.1. *M-sparse solutions of $\mathcal{A}\mathbf{x} = \mathbf{b}$ are unique, if*

$$948 \quad (\text{A.1}) \quad |\langle \mathbf{a}_i, \mathbf{a}_j \rangle| < \frac{1}{2M}, \quad \forall i \neq j,$$

949 *where we assume that the columns of matrix \mathcal{A} are normalized so that $\forall i, \|\mathbf{a}_i\|_{\ell_2} = 1$.*

950 *Proof.* Assume that there exist two M -sparse solutions \mathbf{x}_1 and \mathbf{x}_2 of $\mathcal{A}\mathbf{x} = \mathbf{b}$.
951 Then their difference $\mathbf{z} = \mathbf{x}_1 - \mathbf{x}_2$ is at most $2M$ -sparse, and \mathbf{z} is in the kernel:
952 $\mathcal{A}\mathbf{z} = 0$. This implies that there exist a 1-sparse vector \mathbf{z}_1 and a $(2M - 1)$ -sparse
953 vector \mathbf{z}_2 with disjoint support such that $\mathbf{z}_1 - \mathbf{z}_2 = \mathbf{z}$, and

$$954 \quad (\text{A.2}) \quad \|\mathbf{z}_1\|_{\ell^\infty} \geq \|\mathbf{z}_2\|_{\ell^\infty}.$$

This means that the vector \mathbf{z}_1 was constructed so as to contain only the largest in
magnitude component of \mathbf{z} (one of them if there are several) while \mathbf{z}_2 contains all
the other components of \mathbf{z} . Suppose that the unique non-zero coordinate of \mathbf{z}_1 is i .
Multiplying the identity $\mathcal{A}\mathbf{z}_1 = \mathcal{A}\mathbf{z}_2$ by \mathbf{a}_i , we get

$$\langle \mathbf{a}_i, \mathcal{A}\mathbf{z}_1 \rangle = \langle \mathbf{a}_i, \mathcal{A}\mathbf{z}_2 \rangle,$$

which reduces to

$$(\mathbf{z}_1)_i = \langle \mathbf{a}_i, \mathcal{A}\mathbf{z}_2 \rangle = \sum_{j=1, j \neq i}^{2M} \langle \mathbf{a}_i, \mathbf{a}_j \rangle (\mathbf{z}_2)_j$$

955 Using now (A.1) we obtain

$$956 \quad \|\mathbf{z}_1\|_{\ell^\infty} < \frac{1}{2M}(2M - 1)\|\mathbf{z}_2\|_{\ell^\infty} < \|\mathbf{z}_2\|_{\ell^\infty},$$

957 which is in contradiction with (A.2). □

958 THEOREM 2.2. M -sparse solutions of $\mathcal{A}\mathbf{x} = \mathbf{b}$ can be found as solutions of

959
$$\min \|\mathbf{y}\|_{\ell_1}, \text{ subject to } \mathcal{A}\mathbf{y} = \mathbf{b},$$

960 if

961
$$|\langle \mathbf{a}_i, \mathbf{a}_j \rangle| < \frac{1}{2M}, \quad \forall i \neq j,$$

962 where we assume that the columns of matrix \mathcal{A} are normalized so that $\forall i, \|\mathbf{a}_i\|_{\ell_2} = 1$.

Proof. Assume that there exist two solutions \mathbf{x}_1 and \mathbf{x}_2 of $\mathcal{A}\mathbf{x} = \mathbf{b}$. Suppose x_1 is M -sparse, and x_2 is arbitrary. Their difference $\mathbf{z} = \mathbf{x}_1 - \mathbf{x}_2$ is in the kernel: $\mathcal{A}\mathbf{z} = 0$. We will show that $\|\mathbf{x}_1\|_{\ell_1} < \|\mathbf{x}_2\|_{\ell_1}$. Without loss of generality, we may assume that \mathbf{x}_1 and \mathbf{x}_2 have disjoint support. Otherwise we decompose \mathbf{z} in \mathbf{z}_1 and \mathbf{z}_2 such that $\mathbf{z} = \mathbf{z}_1 - \mathbf{z}_2$ and

$$\begin{aligned} \text{supp}(\mathbf{z}_1) &\subset \text{supp}(\mathbf{x}_1), \\ \text{supp}(\mathbf{z}_2) \cap \text{supp}(\mathbf{x}_1) &= \emptyset. \end{aligned}$$

963 If we assume

964 (A.3)
$$\|\mathbf{x}_2\|_{\ell_1} < \|\mathbf{x}_1\|_{\ell_1}$$

965 then necessarily

966 (A.4)
$$\|\mathbf{z}_2\|_{\ell_1} < \|\mathbf{z}_1\|_{\ell_1}.$$

967 Indeed, if $\|\mathbf{z}_1\|_{\ell_1} \geq \|\mathbf{x}_1\|_{\ell_1}$, it is obvious that (A.3) implies (A.4). Otherwise, if
968 $\|\mathbf{z}_1\|_{\ell_1} < \|\mathbf{x}_1\|_{\ell_1}$ we have

969
$$\|\mathbf{z}_1 - \mathbf{x}_1\|_{\ell_1} \geq \|\mathbf{x}_1\|_{\ell_1} - \|\mathbf{z}_1\|_{\ell_1} > 0.$$

970 Since $\mathbf{x}_2 = \mathbf{x}_1 - \mathbf{z} = \mathbf{x}_1 - \mathbf{z}_1 + \mathbf{z}_2$ we obtain $\|\mathbf{x}_2\|_{\ell_1} = \|\mathbf{x}_1 - \mathbf{z}_1\|_{\ell_1} + \|\mathbf{z}_2\|_{\ell_1}$ and from
971 (A.3) we get

972
$$\|\mathbf{x}_1\|_{\ell_1} > \|\mathbf{x}_2\|_{\ell_1} = \|\mathbf{x}_1 - \mathbf{z}_1\|_{\ell_1} + \|\mathbf{z}_2\|_{\ell_1},$$

973 which implies

974
$$\|\mathbf{z}_2\|_{\ell_1} < \|\mathbf{x}_1\|_{\ell_1} - \|\mathbf{z}_1 - \mathbf{x}_1\|_{\ell_1} \leq \|\mathbf{z}_1\|_{\ell_1}.$$

975 This finishes the proof of the statement that (A.3) implies (A.4).

976 We return now in the proof of the theorem and let i be the coordinate of the
977 component of $\mathbf{z} = \mathbf{z}_1 - \mathbf{z}_2$ with the largest absolute value. Without loss of generality,
978 we may suppose this component is real and positive. Then by multiplying the identity
979 $\mathcal{A}\mathbf{z} = 0$ by \mathbf{a}_i we conclude

980
$$\|\mathbf{z}\|_{\ell_\infty} \leq \frac{1}{2M} \sum_{j \neq i} |z_j| < \frac{1}{2M} \|\mathbf{z}\|_{\ell_1} = \frac{1}{2M} (\|\mathbf{z}_1\|_{\ell_1} + \|\mathbf{z}_2\|_{\ell_1}).$$

981 Since $\|\mathbf{z}_1\|_{\ell_1} \leq M\|\mathbf{z}_1\|_{\ell_\infty} \leq M\|\mathbf{z}\|_{\ell_\infty}$, we obtain

982
$$\|\mathbf{z}\|_{\ell_\infty} < \frac{1}{2} \|\mathbf{z}\|_{\ell_\infty} + \frac{1}{2M} \|\mathbf{z}_2\|_{\ell_1}. \quad \square$$

983 It implies $M\|\mathbf{z}\|_{\ell_\infty} < \|\mathbf{z}_2\|_{\ell_1}$. Again using $\|\mathbf{z}_1\|_{\ell_1} \leq M\|\mathbf{z}\|_{\ell_\infty}$, we obtain $\|\mathbf{z}_1\|_{\ell_1} <$
984 $\|\mathbf{z}_2\|_{\ell_1}$ which is in contradiction with (A.4).

985 THEOREM 2.3. Let \mathbf{x} be a solution of $\mathcal{A}\mathbf{x} = \mathbf{b}$. Let T be the index set of the
 986 support of \mathbf{x} :

$$987 \quad T = \text{supp}(\mathbf{x}), \quad M = |T|.$$

988 Fix a positive $\varepsilon < 1/2$ and suppose that \mathcal{A} satisfies

- 989 i. The columns of matrix \mathcal{A} are normalized so that $\forall i, \|\mathbf{a}_i\|_{\ell_2} = 1$.
 990 ii. The vectors \mathbf{a}_i in the set T are approximately orthogonal, that is they satisfy

$$991 \quad |\langle \mathbf{a}_i, \mathbf{a}_j \rangle| < \frac{\varepsilon}{M}, \quad \forall i, j \in T, i \neq j.$$

992 iii. For any $j \in T$ the vicinity S_j defined as

$$993 \quad S_j = \left\{ k \neq j \mid |\langle \mathbf{a}_k, \mathbf{a}_j \rangle| \geq \frac{1}{2M} \right\},$$

994 has the properties

$$995 \quad |\langle \mathbf{a}_k, \mathbf{a}_j \rangle| \leq 1 - 2\varepsilon, \quad \forall k \in S_j$$

996 and

$$997 \quad |\langle \mathbf{a}_k, \mathbf{a}_j \rangle| < \frac{\varepsilon}{M}, \quad \forall k \in S_i, \forall i \neq j.$$

998 Then \mathbf{x} , the M -sparse solution of $\mathcal{A}\mathbf{x} = \mathbf{b}$, can be found as the solution of

$$999 \quad \min \|\mathbf{y}\|_{\ell_1}, \quad \text{subject to } \mathcal{A}\mathbf{y} = \mathbf{b}.$$

1000

1001 *Proof.* Assume \mathbf{y} is another solution of $\mathcal{A}\mathbf{x} = \mathbf{b}$. Then $\mathcal{A}\mathbf{x} = \mathcal{A}\mathbf{y}$. As in the proof
 1002 of Theorem 2.2 we may suppose that \mathbf{x} and \mathbf{y} have disjoint support. For any $p \in T$
 1003 multiplying the identity $\mathcal{A}\mathbf{x} = \mathcal{A}\mathbf{y}$ by \mathbf{a}_p we get

$$1004 \quad \begin{aligned} x_p + \sum_{i \in T, i \neq p} \langle \mathbf{a}_i, \mathbf{a}_p \rangle x_i &= \sum_{i \in S_p} \langle \mathbf{a}_i, \mathbf{a}_p \rangle y_i + \sum_{i \notin \cup_j S_j} \langle \mathbf{a}_i, \mathbf{a}_p \rangle y_i + \sum_{i \in S_j, j \neq p} \langle \mathbf{a}_i, \mathbf{a}_p \rangle y_i \\ &\leq (1 - 2\varepsilon) \sum_{i \in S_p} |y_i| + \frac{1}{2M} \sum_{i \notin \cup_j S_j} |y_i| + \frac{\varepsilon}{M} \sum_{i \in S_j, j \neq p} |y_i|. \end{aligned}$$

1005 This implies

$$1006 \quad |x_p| < (1 - 2\varepsilon) \sum_{i \in S_p} |y_i| + \frac{1}{2M} \sum_{i \notin \cup_j S_j} |y_i| + \frac{\varepsilon}{M} \sum_{i \in S_j, j \neq p} |y_i| + \frac{\varepsilon}{M} \|\mathbf{x}\|_{\ell_1}.$$

1007 Adding up the inequalities for all $p \in T$ we obtain

$$1008 \quad \|\mathbf{x}\|_{\ell_1} < (1 - \varepsilon) \sum_{i \in \cup_j S_j} |y_i| + \varepsilon \|\mathbf{x}\|_{\ell_1} + \frac{1}{2} \sum_{i \notin \cup_j S_j} |y_i|.$$

1009 Thus

$$1010 \quad (\text{A.5}) \quad \|\mathbf{x}\|_{\ell_1} < \sum_{i \in \cup_j S_j} |y_i| + \frac{1}{2(1 - \varepsilon)} \sum_{i \notin \cup_j S_j} |y_i| \leq \|\mathbf{y}\|_{\ell_1}.$$

1011 Contradiction. □

THEOREM 2.4. **Noisy case** Let \mathbf{x} be an M -sparse solution of

$$\mathcal{A}\mathbf{x} = \mathbf{b},$$

1012 and let as before T denote the index set of the support of \mathbf{x} , that is $T = \text{supp}(\mathbf{x})$ and
 1013 $M = |T|$. Fix a positive $\varepsilon < 1/2$ and suppose that \mathcal{A} satisfies conditions *i*, *ii*, and *iii*
 1014 of Theorem 2.3.

1015 Furthermore, let \mathbf{x}_δ be the ℓ_1 -norm minimal solution of the noisy problem

$$1016 \quad (\text{A.6}) \quad \min \|\mathbf{y}\|_{\ell_1}, \text{ subject to } \mathcal{A}\mathbf{y} = \mathbf{b}^\delta,$$

1017 with \mathbf{b}^δ defined by

$$1018 \quad \mathbf{b}^\delta = \mathbf{b} + \delta\mathbf{b}.$$

We assume that the noise $\delta\mathbf{b}$ is bounded, that is we have

$$\|\delta\mathbf{b}\|_{\ell_2} \leq \delta,$$

1019 for some small positive δ . We further assume that \mathcal{A} has the property that the solution
 1020 $\delta\mathbf{x}$ of

$$1021 \quad (\text{A.7}) \quad \min \|\mathbf{y}\|_{\ell_1}, \text{ subject to } \mathcal{A}\mathbf{y} = \delta\mathbf{b},$$

1022 satisfies

$$1023 \quad (\text{A.8}) \quad \|\delta\mathbf{x}\|_{\ell_1} \leq C\|\delta\mathbf{b}\|_{\ell_2}.$$

1024 Then we can show that the solution \mathbf{x}_δ of (A.6) can be decomposed as

$$1025 \quad (\text{A.9}) \quad \mathbf{x}_\delta = \mathbf{x}_c + \mathbf{x}_i,$$

1026 with \mathbf{x}_c the coherent part of the solution that is supported on T or in the vicinities S_j
 1027 with $j \in T$, and \mathbf{x}_i the incoherent part of the solution which is supported away from
 1028 the vicinities and is small. Specifically, for \mathbf{x}_c we have: for any $j \in T$

$$1029 \quad \left| |(\mathbf{x})_j| - |(\mathbf{x}_c)_j| + \sum_{k \in S_j} \langle a_j, a_k \rangle (\mathbf{x}_c)_k \right| \leq \delta_0 + C\delta,$$

1030 with

$$1031 \quad \delta_0 = \frac{2C\delta(1-\varepsilon)}{M(1-2\varepsilon)} + \frac{2\varepsilon(\|\mathbf{x}\|_{\ell_1} + C\delta)}{M}.$$

While for \mathbf{x}_i we can show that:

$$\|\mathbf{x}_i\|_{\ell_1} \leq \delta_1,$$

1032 with δ_1 given by

$$1033 \quad \delta_1 = C\delta + \frac{4C\delta(1-\varepsilon)}{(1-2\varepsilon)}$$

1034

1035 *Proof.* By assumption (A.7)-(A.8) there exist $\delta\mathbf{x}$ such that $\mathcal{A}\delta\mathbf{x} = \delta\mathbf{b}$, and
 1036 $\|\delta\mathbf{x}\|_{\ell_1} \leq C\delta$. Suppose \mathbf{x} is the M -sparse solution of $\mathcal{A}\mathbf{x} = \mathbf{b}$. Note that

$$1037 \quad \mathcal{A}(\mathbf{x}_\delta - \delta\mathbf{x}) = \mathbf{b}, \quad \mathcal{A}(\mathbf{x} + \delta\mathbf{x}) = \mathbf{b}^\delta.$$

1038 Since both \mathbf{x} and \mathbf{x}_δ are respective minimizers, we obtain

1039 (A.10)
$$\|\mathbf{x}\|_{\ell_1} \leq \|\mathbf{x}_\delta - \delta\mathbf{x}\|_{\ell_1},$$

1040 and

1041
$$\|\mathbf{x}_\delta\|_{\ell_1} \leq \|\mathbf{x} + \delta\mathbf{x}\|_{\ell_1}.$$

1042 Using the triangle inequalities

1043
$$\|\mathbf{x}_\delta - \delta\mathbf{x}\|_{\ell_1} \leq \|\mathbf{x}_\delta\|_{\ell_1} + \|\delta\mathbf{x}\|_{\ell_1}, \quad \|\mathbf{x} + \delta\mathbf{x}\|_{\ell_1} \leq \|\mathbf{x}\|_{\ell_1} + \|\delta\mathbf{x}\|_{\ell_1}$$

1044 we obtain

1045
$$\|\mathbf{x}_\delta - \delta\mathbf{x}\|_{\ell_1} \leq \|\mathbf{x}_\delta\|_{\ell_1} + \|\delta\mathbf{x}\|_{\ell_1} \leq \|\mathbf{x} + \delta\mathbf{x}\|_{\ell_1} + \|\delta\mathbf{x}\|_{\ell_1} \leq \|\mathbf{x}\|_{\ell_1} + 2\|\delta\mathbf{x}\|_{\ell_1}$$

1046 which implies

1047 (A.11)
$$\|\mathbf{x}_\delta - \delta\mathbf{x}\|_{\ell_1} \leq \|\mathbf{x}\|_{\ell_1} + 2C\delta.$$

1048 Combining (A.10) and (A.11) we conclude that

1049 (A.12)
$$\|\mathbf{x}\|_{\ell_1} \leq \|\mathbf{x}_\delta - \delta\mathbf{x}\|_{\ell_1} \leq \|\mathbf{x}\|_{\ell_1} + 2C\delta.$$

For any $p \in T$, taking the inner product of

$$\mathcal{A}(\mathbf{x} - \mathbf{x}_\delta + \delta\mathbf{x}) = 0$$

1050 with \mathbf{a}_p we get

1051 (A.13)
$$\begin{aligned} & (\mathbf{x} - \mathbf{x}_\delta + \delta\mathbf{x})_p + \sum_{k \in T, k \neq p} \langle \mathbf{a}_k, \mathbf{a}_p \rangle (\mathbf{x} - \mathbf{x}_\delta + \delta\mathbf{x})_k + \sum_{k \in S_p} \langle \mathbf{a}_k, \mathbf{a}_p \rangle (\delta\mathbf{x} - \mathbf{x}_\delta)_k \\ & + \sum_{k \in S_j, j \neq p} \langle \mathbf{a}_k, \mathbf{a}_p \rangle (\delta\mathbf{x} - \mathbf{x}_\delta)_k - \sum_{k \notin \cup S_j, k \notin T} \langle \mathbf{a}_k, \mathbf{a}_p \rangle (\delta\mathbf{x} - \mathbf{x}_\delta)_k = 0. \end{aligned}$$

1052 Using properties (ii)-(iii) we obtain

1053 (A.14)
$$\begin{aligned} |(\mathbf{x} - \mathbf{x}_\delta + \delta\mathbf{x})_p| & < \frac{\varepsilon}{M} \sum_{k \in T, k \neq p} |(\mathbf{x} - \mathbf{x}_\delta + \delta\mathbf{x})_k| \\ & + (1 - 2\varepsilon) \sum_{k \in S_p} |(\mathbf{x}_\delta - \delta\mathbf{x})_k| + \frac{\varepsilon}{M} \sum_{k \in S_j, j \neq p} |(\mathbf{x}_\delta - \delta\mathbf{x})_k| \\ & + \frac{1}{2M} \sum_{k \notin \cup S_j, k \notin T} |(\mathbf{x}_i - \delta\mathbf{x})_k|. \end{aligned}$$

1054 Summing over all $p \in T$ we get

1055
$$\begin{aligned} \sum_{p \in T} |(\mathbf{x} - \mathbf{x}_\delta + \delta\mathbf{x})_p| & < \varepsilon \sum_{p \in T} |(\mathbf{x} - \mathbf{x}_\delta + \delta\mathbf{x})_p| + (1 - 2\varepsilon) \sum_{k \in \cup_{p=1}^M S_p} |(\mathbf{x}_\delta - \delta\mathbf{x})_k| \\ & + \varepsilon \sum_{k \in \cup_{p=1}^M S_p} |(\mathbf{x}_\delta - \delta\mathbf{x})_k| + \frac{1}{2} \sum_{k \notin \cup S_j, k \notin T} |(\mathbf{x}_i - \delta\mathbf{x})_k|. \end{aligned}$$

1056 Thus

1057
$$\begin{aligned} \sum_{k \in T} |(\mathbf{x} - \mathbf{x}_\delta + \delta\mathbf{x})_k| & < \sum_{k \in \cup_{p=1}^M S_p} |(\mathbf{x}_\delta - \delta\mathbf{x})_k| + \frac{1}{2(1 - \varepsilon)} \sum_{k \notin \cup S_j, k \notin T} |(\mathbf{x}_i - \delta\mathbf{x})_k| \\ & = \sum_{k \notin T} |(\mathbf{x}_\delta - \delta\mathbf{x})_k| - \frac{1 - 2\varepsilon}{2(1 - \varepsilon)} \sum_{k \notin \cup S_j, k \notin T} |(\mathbf{x}_i - \delta\mathbf{x})_k| \end{aligned}$$

1058 We therefore obtain

$$1059 \quad \|\mathbf{x}\|_{\ell_1} < \|\mathbf{x}_\delta - \delta\mathbf{x}\|_{\ell_1} - \frac{1-2\varepsilon}{2(1-\varepsilon)} \sum_{k \notin \cup S_j, k \notin T} |(\mathbf{x}_i - \delta\mathbf{x})_k|$$

1060 By (A.12) we conclude

$$1061 \quad \sum_{k \notin \cup S_j, k \notin T} |(\mathbf{x}_i - \delta\mathbf{x})_k| \leq \frac{4C\delta(1-\varepsilon)}{1-2\varepsilon}.$$

1062 By the triangle inequality

$$1063 \quad (\text{A.15}) \quad \|\mathbf{x}_i\|_{\ell_1} \leq \|\delta\mathbf{x}\|_{\ell_1} + \frac{4C\delta(1-\varepsilon)}{1-2\varepsilon} \leq C\delta + \frac{4C\delta(1-\varepsilon)}{1-2\varepsilon} = \delta_1.$$

1064 It remains to investigate \mathbf{x}_c , the coherent part of the solution. From (A.13) we have

$$\begin{aligned} \left| (\mathbf{x})_p + \sum_{k \in S_p \cup \{p\}} \langle \mathbf{a}_k, \mathbf{a}_p \rangle (\delta\mathbf{x} - \mathbf{x}_\delta)_k \right| &< \frac{\varepsilon}{M} \sum_{k \in T, k \neq p} |(\mathbf{x} - \mathbf{x}_\delta + \delta\mathbf{x})_k| + \frac{\varepsilon}{M} \sum_{k \in S_j, j \neq p} |(\mathbf{x}_\delta - \delta\mathbf{x})_k| \\ &+ \frac{1}{2M} \sum_{k \notin \cup S_j, k \notin T} |(\mathbf{x}_i - \delta\mathbf{x})_k| \\ 1065 \quad &\leq \frac{\varepsilon}{M} \|\mathbf{x} - \mathbf{x}_\delta + \delta\mathbf{x}\|_{\ell_1} + \frac{1}{2M} \frac{4C\delta(1-\varepsilon)}{1-2\varepsilon} \\ &\leq \frac{\varepsilon}{M} (\|\mathbf{x}\|_{\ell_1} + \|\mathbf{x}_\delta\|_{\ell_1} + \|\delta\mathbf{x}\|_{\ell_1}) + \frac{2C\delta(1-\varepsilon)}{M(1-2\varepsilon)} \\ &\leq \frac{\varepsilon}{M} (2\|\mathbf{x}\|_{\ell_1} + 2C\delta) + \frac{2C\delta(1-\varepsilon)}{M(1-2\varepsilon)} = \delta_0. \end{aligned}$$

1066 Applying the triangle inequality:

$$\begin{aligned} 1067 \quad \left| (\mathbf{x})_p - \sum_{k \in S_p \cup \{p\}} \langle \mathbf{a}_k, \mathbf{a}_p \rangle (\mathbf{x}_\delta)_k \right| &\leq \left| (\mathbf{x})_p + \sum_{k \in S_p \cup \{p\}} \langle \mathbf{a}_k, \mathbf{a}_p \rangle (\delta\mathbf{x} - \mathbf{x}_\delta)_k \right| + \left| \sum_{k \in S_p} \langle \mathbf{a}_k, \mathbf{a}_p \rangle (\delta\mathbf{x})_k \right| \\ 1068 \quad &\leq \delta_0 + C\delta, \\ 1069 \end{aligned}$$

1070 we obtain the result. \square

1071 **Appendix B. Proof of theorem 2.6.**

1072 **THEOREM 2.6.** Let $X = \text{Diag}(\mathbf{x})$ be a diagonal matrix that solves

$$1073 \quad \tilde{\mathcal{A}}XL = B,$$

1074 where $\tilde{\mathcal{A}}$ satisfies conditions (i), (ii), and (iii) of Theorem 2.3 for a fixed $\varepsilon < 1/3$,

$$1075 \quad L = \begin{pmatrix} l_{11} & l_{12} & l_{1S} \\ l_{21} & l_{22} & l_{2S} \\ \vdots & \vdots & \vdots \\ l_{K1} & l_{K2} & l_{KS} \end{pmatrix} \in \mathbb{C}^{K \times S},$$

1076 and B is the noiseless data matrix (2.17) with SVD $B = Q = U\Sigma V^T$. Let the
1077 perturbed matrix $B^\delta = Q^\delta + Q_0$ be such that $\sigma_{\max}(B^\delta - B) \leq \delta$. Suppose \mathbf{x} , the
1078 vector diagonal entries of X , is sparse with $T = \text{supp}(\mathbf{x})$, $M = |T|$, $M \ll \text{size}(\mathbf{x})$, and

$$1079 \quad x_m = \min_{x_i \neq 0} \{|x_i|\}.$$

1080 Let L_T be the submatrix of L , formed by the rows corresponding to T , has

$$1081 \quad (\text{B.1}) \quad \sigma_m^T = \sigma_{\min}(L_T).$$

1082 If

$$1083 \quad (\text{B.2}) \quad 2\delta < x_m \sigma_m^T (1 - 3\varepsilon),$$

1084 the orthogonal projections onto the subspaces $R(Q^\delta)$ and $R(B)$ are close:

$$1085 \quad (\text{B.3}) \quad \|P_{R(Q^\delta)} - P_{R(B)}\|_{\ell_2} \leq \frac{\delta}{x_m \sigma_m^T (1 - 3\varepsilon)}.$$

1086 *Proof.* Denote by X_T be the submatrix of X where we keep the rows that corre-
1087 spond to the support of \mathbf{x} . Similarly, denote by \mathbf{y}_T be the subvector of \mathbf{y} where we
1088 keep the entries that correspond to the support of \mathbf{x} . We claim that

$$1089 \quad (\text{B.4}) \quad (1 - 3\varepsilon)^2 \|\mathbf{z}\|_{\ell_2}^2 \leq \|(\tilde{\mathcal{A}}^* \mathbf{z})_T\|_{\ell_2}^2 \leq (1 + 3\varepsilon)^2 \|\mathbf{z}\|_{\ell_2}^2$$

1090 if $\mathbf{z} \in R(B)$. Indeed, suppose that

$$1091 \quad \mathbf{z} = \sum_{i \in T} \alpha_i \mathbf{a}_i.$$

1092 Then, defining $\boldsymbol{\alpha}$ as the vector in \mathbb{C}^K whose components are zero except the i th
1093 components with $i \in T$ that are equal to α_i , we get

$$1094 \quad \left| \|\mathbf{z}\|_{\ell_2}^2 - \|\boldsymbol{\alpha}\|_{\ell_2}^2 \right| = \left| \sum_{i,j \in T, i \neq j} \bar{\alpha}_i \alpha_j \langle \mathbf{a}_i, \mathbf{a}_j \rangle \right| \leq \varepsilon \|\boldsymbol{\alpha}\|_{\ell_2}^2,$$

1095 and

$$1096 \quad (1 - \varepsilon) \|\boldsymbol{\alpha}\|_{\ell_2}^2 \leq \|\mathbf{z}\|_{\ell_2}^2 \leq (1 + \varepsilon) \|\boldsymbol{\alpha}\|_{\ell_2}^2.$$

1097 For any $j \in T$ we have

$$1098 \quad (\tilde{\mathcal{A}}^* \mathbf{z})_j = \sum_{i \in T} \alpha_i \langle \mathbf{a}_j, \mathbf{a}_i \rangle,$$

1099 and, therefore,

$$1100 \quad \|(\tilde{\mathcal{A}}^* \mathbf{z})_T\|_{\ell_2}^2 = \sum_{i,j,k \in T} \bar{\alpha}_j \alpha_i \langle \mathbf{a}_k, \mathbf{a}_i \rangle \overline{\langle \mathbf{a}_k, \mathbf{a}_j \rangle}.$$

1101 Hence,

$$1102 \quad \left| \|(\tilde{\mathcal{A}}^* \mathbf{z})_T\|_{\ell_2}^2 - \|\boldsymbol{\alpha}\|_{\ell_2}^2 \right| \leq \left| \sum_{i,j,k \in T, i \neq j} \bar{\alpha}_j \alpha_i \langle \mathbf{a}_k, \mathbf{a}_i \rangle \overline{\langle \mathbf{a}_k, \mathbf{a}_j \rangle} \right|$$

$$1103 \quad \leq \sum_{i,j \in T, i \neq j} \frac{|\alpha_j|^2 + |\alpha_i|^2}{2} \varepsilon \left(\frac{2}{M} + \frac{\varepsilon}{M} \right) \leq 3\varepsilon \|\boldsymbol{\alpha}\|_{\ell_2}^2.$$

1105 Therefore,

$$1106 \quad (1 - 3\varepsilon) \|\boldsymbol{\alpha}\|_{\ell_2}^2 \leq \|(\tilde{\mathcal{A}}^* \mathbf{z})_T\|_{\ell_2}^2 \leq (1 + 3\varepsilon) \|\boldsymbol{\alpha}\|_{\ell_2}^2,$$

1107 and we obtain

$$1108 \quad \frac{1 - 3\varepsilon}{1 + \varepsilon} \|\mathbf{z}\|_{\ell_2}^2 \leq \|(\tilde{\mathcal{A}}^* \mathbf{z})_T\|_{\ell_2}^2 \leq \frac{1 + 3\varepsilon}{1 - \varepsilon} \|\mathbf{z}\|_{\ell_2}^2,$$

1109 which implies (B.4).

1110 In order to compute the smallest nonzero singular value of B we observe that

$$\begin{aligned}
1111 \quad & \min_{\mathbf{z} \in R(B), \|\mathbf{z}\|_{\ell_2}=1} \mathbf{z}^* B B^* \mathbf{z} = \min_{\mathbf{z} \in R(\tilde{B}), \|\mathbf{z}\|_{\ell_2}=1} (\tilde{\mathcal{A}}^* \mathbf{z})_T^* X_T L_T L_T^* \bar{X}_T (\tilde{\mathcal{A}}^* \mathbf{z})_T \\
1112 \quad & \geq (1 - 3\varepsilon)^2 \min_{\mathbf{y} \in \mathbb{C}^M, \|\mathbf{y}\|_{\ell_2}=1} \mathbf{y}^* X_T L_T L_T^* \bar{X}_T \mathbf{y} \geq (1 - 3\varepsilon)^2 x_m^2 (\sigma_m^T)^2, \\
1113
\end{aligned}$$

1114 where we have used the condition (B.1). Since $\sigma_{\max}(B^\delta - B) \leq \delta$, we conclude that
1115 $B^\delta = Q^\delta + Q_0^\delta$, where Q^δ has M nonzero singular values, with smallest nonzero
1116 singular value

$$1117 \quad \sigma_{\min}(Q^\delta) \geq x_m \sigma_m^T (1 - 3\varepsilon) - \delta,$$

1118 and Q_0^δ has largest singular value

$$1119 \quad \sigma_{\max}(Q_0^\delta) \leq \delta.$$

1120 If (B.2) holds, then we can discard Q_0^δ by truncation of the singular values smaller
1121 than the noise level. We now apply Theorem 2.5 to obtain (B.3). \square